

Electrical properties of (11-22) Si:AlGa_xN layers at high Al contents grown by metal-organic vapor phase epitaxy

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Abstract

In this work, the growth and conductivity of semipolar Al_xGa_{1-x}N:Si with (11-22) orientation is investigated. Al_xGa_{1-x}N:Si ($x = 0.60 \pm 0.03$ and $x = 0.80 \pm 0.02$) layers were grown with different SiH₄ partial pressures and the electrical properties were determined using Hall measurements at room temperature. The aluminum mole fraction was measured by wavelength dispersive X-ray spectroscopy and X-ray diffraction and the Si-concentration was measured by wavelength dispersive X-ray spectroscopy and secondary ion mass spectroscopy. Layer resistivities as low as 0.024 Ω cm for $x = 0.6$ and 0.042 Ω cm for $x = 0.8$ were achieved. For both aluminum mole fractions the resistivity exhibits a minimum with increasing Si concentration which can be explained by compensation due to formation of cation vacancy complexes at high doping levels. The onset of self-compensation occurs at larger estimated Si concentrations for larger Al contents.

The efficiency of UVC-LEDs emitting below 240 nm drops strongly with decreasing emission wavelength [1, 2]. One of the reasons is that LEDs grown on c-plane sapphire suffer from a change in the optical polarization from dominant transverse electric (TE) polarized emission to dominant transverse magnetic (TM) polarized emission at 240 nm [3, 4], which subsequently leads to a reduced light extraction efficiency (LEE) [5-7]. However, UVC-LEDs on semipolar planes, i.e. (11-22), overcome poor LEE through a rotation of the wurtzite unit cell [8]. Furthermore, semipolar MQWs exhibit reduced internal polarization fields, which leads to a high electron and hole wave function overlap and high radiative recombination rate [6, 9, 10]. The (11-22) orientation can be nucleated on (10-10) m-plane sapphire [11], and has shown success in InGaN devices [12, 13]. However, there has only been one report of a UV LED consisting of (11-22) AlGaN layers, emitting at 307 nm [14]

To realize an efficient semipolar UV LED, highly conductive n-doped layers are essential as current is transported laterally from the n-contacts to the pn-junction and high layer resistance leads to current crowding [15]. However, there is very little literature on Si-doping of (11-22) AlGaN layers [16, 17], especially at higher Al contents. In this work we present the growth of conductive (11-22) AlGaN:Si on (10-10) m-plane sapphire by metal organic vapor phase epitaxy (MOVPE) at an Al contents of 60% and 80%. Such n-layers would be suitable for the realization of LEDs emitting at 270 nm and 240 nm, respectively [18]. The Si-concentration and AlN mole fraction were determined by wavelength dispersive x-ray spectroscopy (WDX) and the electrical properties were determined by room temperature Hall measurements.

MOVPE growth was performed in a closed coupled showerhead reactor using trimethylaluminum (TMAI), trimethylgallium (TMGa), silane (SiH₄), and ammonia (NH₃) on 2''

on-axis m-plane (10-10) sapphire substrates. The sapphire substrates underwent 600 seconds of nitridation at a temperature of 1135 °C, at a reactor pressure of 500 hPa, and with an NH₃ partial pressure of 30 Pa.

The growth of AlN and unintentionally doped AlGa_N used in the structures was tailored to mitigate and eliminate the density of misoriented grains on the surface and are detailed elsewhere [19]. For the samples grown at 60% Al, the TMAI/ (TMGa+TMAI) partial pressure ratio in the gas phase was 0.60 and the V/III ratio was 470. The carrier gas flow into the reactor was 20 standard liters per minute (slm) of hydrogen at a reactor pressure of 500 hPa. The growth temperature was 1095 °C. The TMAI, TMGa, and NH₃ partial pressures were 5 Pa, 3 Pa, and 3750 Pa, respectively. The SiH₄/III ratio was varied from 1.5x10⁻⁵ to 3x10⁻⁴. For the samples grown at 80% Al, the TMAI/ (TMGa+TMAI) partial pressure ratio in the gas phase was 0.76. The mean growth temperature was 1100 °C and a V/III ratio of 630. The TMAI, TMGa, and NH₃ partial pressures were 5 Pa, 1 Pa, and 3750 Pa, respectively. The SiH₄/III ratio was varied from 5x10⁻⁵ to 6x10⁻⁴. The thickness of the AlGa_N layers was determined using in-situ spectroscopic reflectance using a Laytec EpiCurve TT system. At 60% Al, the film thickness of the layer was on average 1.15 μm. At 80% Al, the film thickness of the layer was on average 0.55 μm. The (11-22) growth orientation in the AlGa_N layers was confirmed by x-ray diffraction (XRD) for each sample with an omega-two theta scan using an open detector in double axis configuration with a four-bounce monochromator in the incident beam. No additional peaks, such as the (1-10-3) AlGa_N, which has been the subject of discussion in other works [20-22], were observed. A representative ω/2θ scan exhibiting the sapphire (30-30) and the AlN and AlGa_N (11-22) reflexes is shown in Figure 1.a.

The crystal quality was assessed by XRD open detector rocking curve scans of the

symmetric (11-22) reflection. All samples exhibited a FWHM of 1800'' with incidence along the [10-10] and 900'' with incidence along the [11-2-3] in-plane direction for all samples hinting to a high density of threading dislocations in the range of $10^{10} - 10^{11} \text{ cm}^{-2}$ [23, 24]. The FWHM of the omega rocking curves scans showed no dependence on the Si concentration or Al content and remained constant throughout all samples. The surface morphology and root-mean squared (RMS) roughness were investigated and determined using an atomic force microscope (AFM) in contact mode. As the silicon concentration is increased, the surface morphology and roughness remained unchanged and is shown to be independent of the increasing silicon concentration. Fig.1b shows a representative AFM image of the surface morphology which an RMS roughness of 3.08 nm at a SiH_4/III of 1.25×10^4 at an Al mole fraction of 60%.

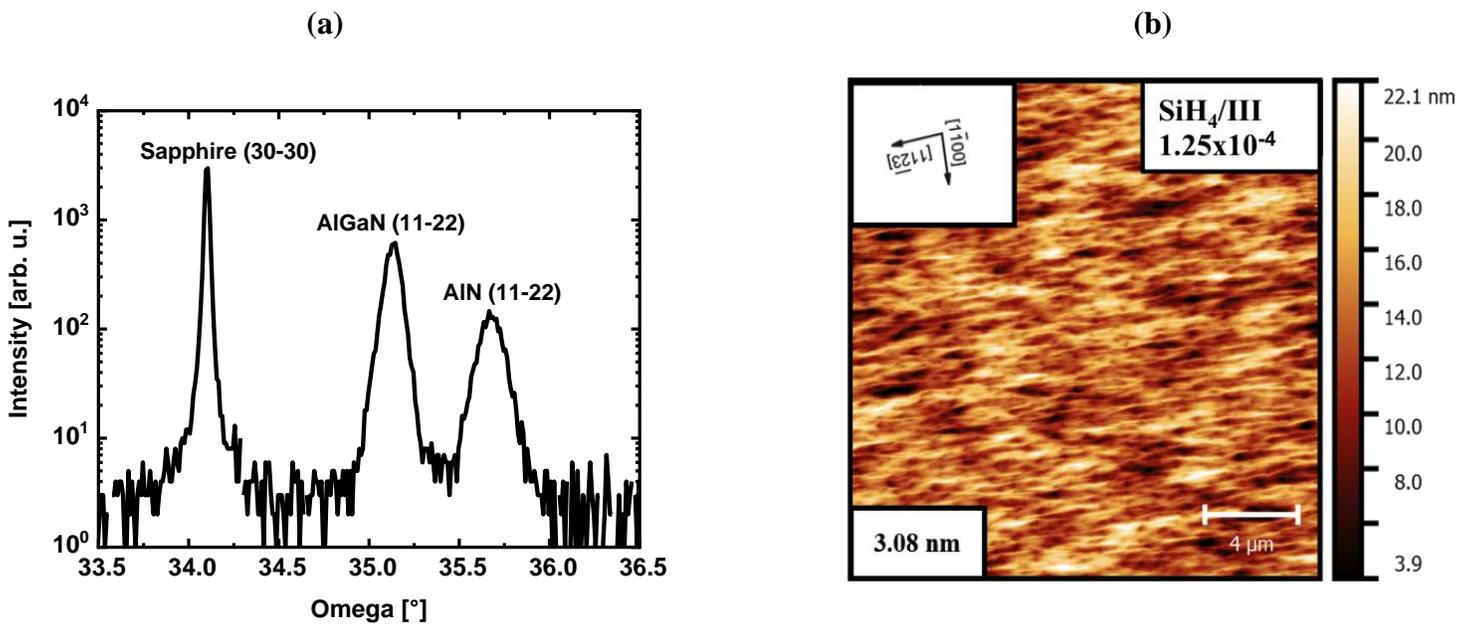


Figure 1 (a) Symmetric HR-XRD $\omega/2\theta$ scan of AlGaN:Si / AlN on m-plane sapphire. (b) $20 \mu\text{m} \times 20 \mu\text{m}$ AFM images of the n-AlGaN surface at SiH_4/III of 1.25×10^4 with an RMS roughness of 3.08 nm.

The WDX measurements in this work were performed on an electron probe microanalyser (JEOL JXA-8530F) using a 10 keV electron beam. Monte Carlo simulations, using CASINO software [25] indicate that this beam energy corresponds to deposition of 90% of the beam energy within an approximate depth of 500 nm for high AlN mole fraction ($\text{Al}_{0.7}\text{Ga}_{0.3}\text{N}$) and 450nm for the higher density lower AlN mole fractions ($\text{Al}_{0.5}\text{Ga}_{0.5}\text{N}$).

The WDX measurements to determine the Al, Ga and N compositions were performed using a 40 nA beam current and X-ray counting times of 60 s (for the peak) and 30 s (for the backgrounds). K-Ratios (X-ray intensity from sample divided by that from the standard) were determined by comparing data from the sample to a AlN:Si standard for the Al and N composition and GaN for Ga. Due to the smaller Si concentration the beam current was increased to 400nA with extended counting times (360/180s). Pure Si was used as the standard. Results from both measurements were then combined within the WDX software. A large thallium acid phthalate (TAP) crystal was used to record the Al K_{α} , Ga L_{α} and Si K_{α} X-ray intensities while a layered pseudo-crystal recorded the N K_{α} x-ray intensity [26].

For each sample, WDX analyses were performed across 9 points corresponding to a 3x3 grid. The measured Si concentration was then calibrated using measured secondary ion mass spectroscopy (SIMS) data to provide better estimates of the dopant concentration [27,28]. Random errors were derived for each set of measurements. A systematic error of 2% was assumed for WDX measurements of Al, Ga and N, increasing to 3% for the more challenging measurement of Si. The AlN mole fraction of the AlGaN:Si layers was determined to be on average $60 \pm 3 \%$ and $80 \pm 2 \%$ for the samples measured by WDX.

For Hall effect measurements on $0.8 \times 0.8 \text{ cm}^2$ sample pieces, V/Al/Ni/Au (15/90/20/30 nm) n-contacts were thermally evaporated, followed by a rapid annealing step of 40 s at 800°C

under N₂ atmosphere [29]. The Hall effect was measured in the van-der-Pauw configuration (Lake Shore HMS 7504) with an excitation current, magnetic field, and dwell time of 10 μA, 3 kG, and 2s, respectively. The resistivity was additionally determined by contactless eddy current resistivity measurements (CRM).

Figure 2.a displays the silicon and carrier concentration (at 300 K) as function of the SiH₄/III ratio. The open red and blue points represent the Si concentration as measured by WDX and SIMS in the AlGa_N:Si layers at 60% Al and 80% Al, respectively. The measurements show a linear increase of the Si with the SiH₄ partial pressure in the range of $2.6 \times 10^{18} \text{ cm}^{-3}$ for the lowest SiH₄/III ratio to $2 \times 10^{19} \text{ cm}^{-3}$ for the highest SiH₄/III ratio of 6×10^{-4} .

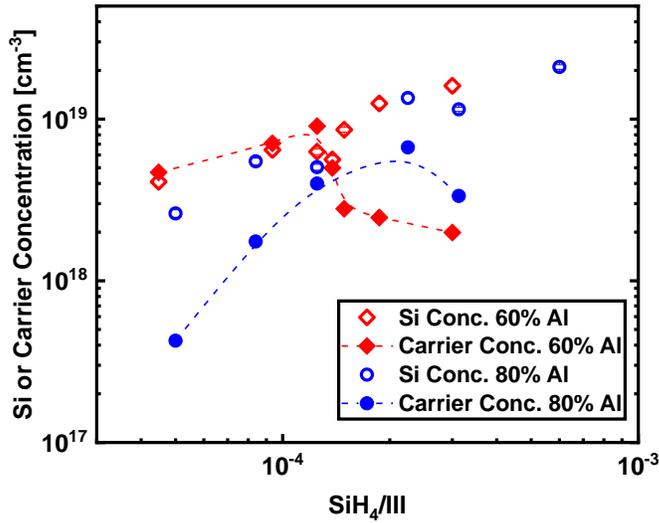
In order to later analyze the electrical properties with respect to the Si concentration in the layers, the linear regression model was applied to Si concentration measured by WDX in the AlGa_N:Si layers at 60% Al and 80% Al, where the sample grown at a SiH₄/III ratio of 6×10^{-4} was excluded for the sake of accuracy. For both aluminum mole fractions the Si incorporation is identical within the scattering of the measured data points. This is consistent with a linear Si incorporation independent of the aluminum mole fraction observed for c-plane Al_xGa_{1-x}N [26]. The y intercept (Si concentration) was set to 0, with the assumption that at a SiH₄/III ratio of 0, the Si concentration is also 0 and is shown graphically in Fig 2.b

The closed red and blue points represent the carrier concentration as determined by room temperature (RT) Hall effect measurements in the AlGa_N:Si layers at an aluminum mole fraction of 60% and 80%, respectively. At an aluminum mole fraction of 60%, the RT carrier concentration is initially in agreement with the Si-concentration measured by WDX, and both are increasing with increasing SiH₄/III ratio, reaching a maximum at a SiH₄/III ratio of 1.25×10^{-4} . After this maximum, the carrier concentration drops sharply while the Si concentration continues

to increase. In this regime, the RT carrier density is less than one fifth of Si concentration or less, suggesting a high degree of compensation at the higher donor concentrations investigated. For the samples with an aluminum mole fraction of 80%, the RT carrier concentration first increases, reaches a maximum at a SiH₄/III ratio (2.25×10^{-4}), and then drops off, while the Si concentration continued to increase. The maximum appears at higher SiH₄/III ratio (and Si concentration) compared to the samples with an aluminum mole fraction of 60%. The carrier density is also significantly below the Si-concentration measured by WDX, which can be caused by a higher ionization energy of the donors [30] and a higher degree of compensation in the samples. Nonetheless, both the Si and carrier concentrations are increasing with increasing SiH₄/III ratio.

A closer look at the measured carrier density ($9 \times 10^{18} \text{ cm}^{-3}$) at the SiH₄/III ratio of 1.25×10^{-4} for an aluminum mole fraction of 60% in Fig. 2.a revealed that the carrier density is larger than the WDX measured silicon concentration ([Si] $6.4 \times 10^{18} \text{ cm}^{-3}$). One possible explanation might be a higher donor concentration caused by unintentionally incorporated oxygen, which is known to incorporate effectively on (11-22) surfaces [31]. SIMS results obtained for this sample show an oxygen concentration of $6 \times 10^{18} \text{ cm}^{-3}$, suggesting a contribution to the carrier concentration. Another explanation could be a scattering of the determined Si concentrations by WDX as this particular sample is below the linear regression. We conclude that both effects play a role showing that the calculation of the compensation ratio [32] based on SIMS, WDX and Hall effect measurements must be carefully reviewed.

(a)



(b)

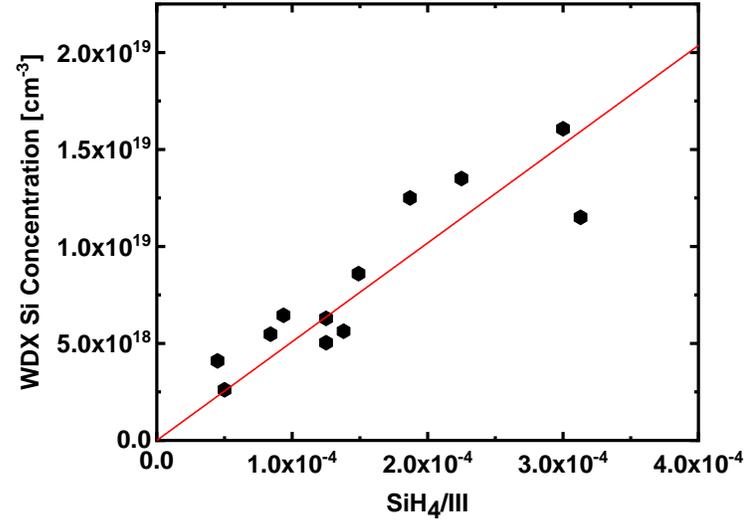


Figure 2 (a) The silicon and carrier concentration determined by WDX and Hall Effect measurement for varying SiH_4/III ratio at 60% Al and 80% Al. The carrier concentrations are marked with guides to the eye **(b)** Linear regression of WDX Si concentration for both Al contents as a function of SiH_4/III .

Figure 3 displays the results for the a. carrier concentration and b. mobility as a function of the SiH_4/III ratio on the lower x-axis and the Si concentration as determined by the linear regression model in the upper x-axis. At an aluminum mole fraction of 60%, an almost constant mobility and an increase in the carrier density with increasing silicon concentration was observed up to a maximum at the estimated Si concentration of $6.4 \times 10^{18} \text{ cm}^{-3}$ (SiH_4/III ratio of 1.25×10^{-4}). The maximum carrier density and mobility reached are $9 \times 10^{18} \text{ cm}^{-3}$ and $24 \text{ cm}^2 \text{V}^{-1} \text{ s}^{-1}$, respectively. Thereafter, both the mobility and carrier density decrease with increasing Si concentration. For the samples with an aluminum mole fraction of 80%, an increase in both the carrier density and mobility with increasing silicon concentration is observed up to a maximum at the estimated Si concentration of $1.1 \times 10^{19} \text{ cm}^{-3}$ (SiH_4/III ratio of 2.25×10^{-4}). At this estimated Si concentration, the carrier density and mobility reached a maximum value of $6.7 \times 10^{18} \text{ cm}^{-3}$ and

21 $\text{cm}^2\text{V}^{-1}\text{s}^{-1}$, respectively. A higher Si concentration yielded a lower carrier concentration and mobility. The sample with the highest SiH_4/III ratio of 6×10^{-4} could not be measured by Hall effect measurements.

This trend is described in the literature as “knee-like behavior”, being observed for AlGaInSi with high aluminum mole fractions [30, 33, 34]. The mobility of 20 - 24 $\text{cm}^2\text{V}^{-1}\text{s}^{-1}$ measured here resembles the value for high dislocation density (\sim mid 10^{10} cm^{-3}) c-plane $\text{Si:Al}_{0.7}\text{Ga}_{0.3}\text{N}$ [33], which indicates that mobility is independent of the growth plane for materials with a high defect density, as also the samples investigated here exhibit a high dislocation density in the range of $10^{10} - 10^{11} \text{ cm}^{-2}$ [23, 24] indicated by the FWHM of the symmetric (11-22) reflex which was equal for all samples of 900” and 1800” with incidence along [11-2-3] and [10-10], respectively.

The subsequent drop in mobility and carrier concentration is attributed to the scattering of ionized impurities along with a high dislocation density [35-37] and to the onset of self-compensation via a cation vacancy complex with Si, as more silane is introduced [38]. This indicates that at higher SiH_4 flows, Si overcompensation occurs in the structures and additional Si leads to the formation of compensating defects instead of a higher net donor density. [39-41].

The “knee-like behavior” occurs at different Si concentrations for the different aluminum mole fractions reported in this work. The carrier density drops at an estimated Si concentration of $1.1 \times 10^{19} \text{ cm}^{-3}$ for an aluminum mole fraction of 80%, compared to an estimated Si concentration of 6.4×10^{18} for an aluminum mole fraction of 60% even though the peak carrier concentration and mobility are slightly lower at an aluminum mole fraction of 80% ($6.7 \times 10^{18} \text{ cm}^{-3}$ and $21 \text{ cm}^2\text{V}^{-1}\text{s}^{-1}$) compared to 60% ($9 \times 10^{18} \text{ cm}^{-3}$ and $24 \text{ cm}^2\text{V}^{-1}\text{s}^{-1}$). This indicates that the critical Si concentration for self-compensation, i.e., where the chemical potential of Si increases and the formation of cation vacancy complexes is favorable, depends on the Al content. This is

in direct contrast to c-plane AlGaIn with aluminum mole fraction in the range of 80% – 96%, where the critical Si concentration decreased with increasing aluminum mole fraction [30, 34], we find here the opposite trend. Further investigations will clarify if this trend can be observed as well for higher and lower aluminum mole fractions and which role oxygen plays [42], as it forms a stable DX center at aluminum mole fractions of 60% [43].

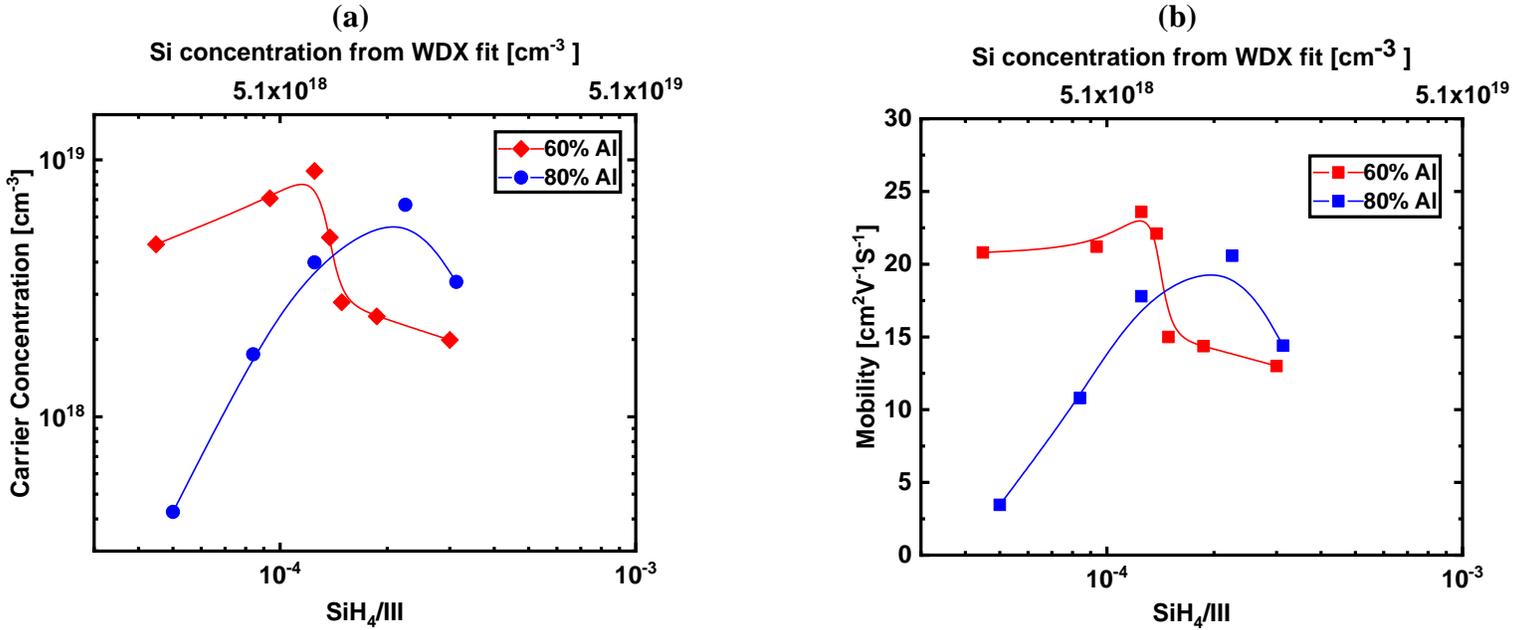


Figure 3 (a) Carrier concentration, and (b) mobility as a function of SiH₄/III ratio and Si concentration fitted from WDX results as measured by RT Hall measurements at 60% Al and 80% Al. The corresponding lines are simply a guide to the eye.

Figure 4 displays the layer resistivity, measured by RT Hall effect measurements and CRM, as a function of the SiH₄/III on the lower x-axis and the Si concentration as determined by the linear regression model in the upper x-axis. For all samples the resistivity values obtained with the two methods are in good agreement for all samples investigated. At an aluminum mole fraction of 60% (red data points) a decrease in the layer resistivity with increasing Si concentration was observed down to a minimum of 0.024 Ω cm, at an estimated Si concentration of 6.4 × 10¹⁸ cm⁻³. Thereafter, the layer resistivity increased till ~ 0.2 Ω cm for the last measured

sample at 60% Al (estimated Si concentration $\sim 1.5 \times 10^{19} \text{ cm}^{-3}$). At an aluminum mole fraction of 80%, a decrease in the layer resistivity with increasing Si concentration was also observed down to a minimum of $0.042 \text{ } \Omega \text{ cm}$, at an estimated Si concentration of $1.1 \times 10^{19} \text{ cm}^{-3}$. The layer resistivity then increased to $\sim 1.1 \text{ } \Omega \text{ cm}$ for the last measured sample at 80% Al (estimated Si concentration $\sim 3 \times 10^{19} \text{ cm}^{-3}$). For both aluminum mole fractions the lowest resistivity is caused by the highest carrier density and mobility as shown in Figure 3. Again Fig. 4 clearly demonstrates that the optimal resistivity value obtained shifts to larger estimated Si concentrations for higher Al contents, the opposite trend of what has observed in the literature for c-plane Al-rich Si:AlGaN [30,34].

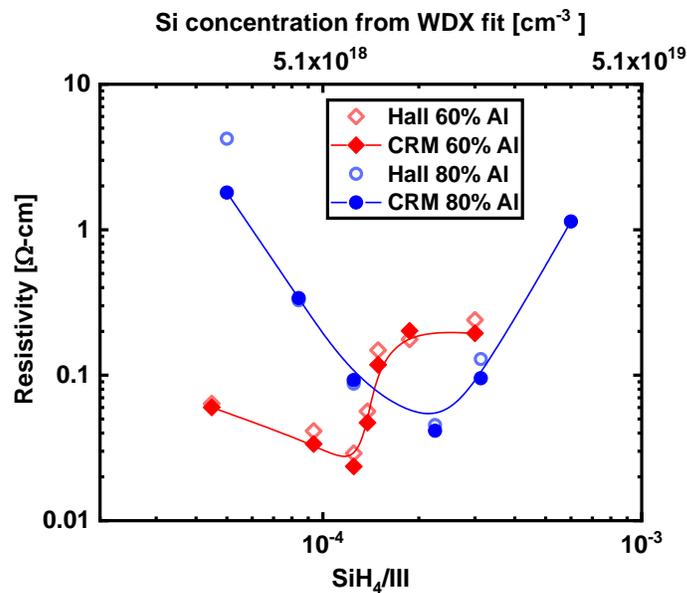


Figure 4 Layer resistivity as a function of SiH₄/III ratio and Si concentration fitted from WDX results as measured by RT Hall measurements and contactless measurements at 60% Al and 80% Al. The corresponding lines are simply a guide to the eye.

In conclusion, the growth and conductivity of (11-22) Al_xGa_{1-x}N:Si ($x = 0.60 \pm 0.03$ and $x = 0.80 \pm 0.02$) at different SiH₄ partial pressures was investigated. The Al content and Si concentration were measured by WDX. The measurements show a linear increase of the Si

concentration with the SiH₄ partial pressure in the range of 2x10¹⁸ cm⁻³ to 3x10¹⁹ cm⁻³. The electrical properties were acquired using contactless resistivity measurements and Hall effect measurements at room temperature. For an aluminum mole fraction of 60% at an estimated Si concentration of 6.4x10¹⁸ cm⁻³, a minimal layer resistivity of 0.024 Ωcm and maximum mobility and carrier density of 24 cm²V⁻¹s⁻¹ and 9x10¹⁸ cm⁻³ were achieved, respectively. For an aluminum mole fraction of 80% Al, an estimated Si concentration of 1.1x10¹⁹ cm⁻³, a minimal layer resistivity of 0.042 Ωcm and maximum mobility and carrier density of 20 cm²V⁻¹s⁻¹ and 6.7x10¹⁸ cm⁻³ were achieved, respectively. Contrary to what is observed in c-plane AlGa_N:Si for higher aluminum mole fractions, the peak electrical properties are observed at larger estimated silicon concentrations for larger Al contents. Therefore, the onset of self-compensation via a cation vacancy complex Si is dependent not only on the Si concentration and dislocation density, but also the Al content.

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Data Availability

The data that support the findings of this study are available from the corresponding author upon reasonable request

References

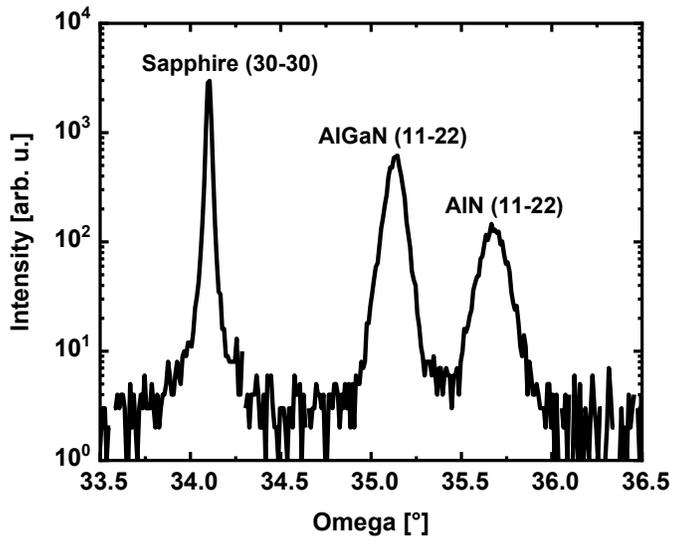
[1] M.Kneissl, J.Rass, III-Nitride Ultraviolet Emitters, Technology and Applications (Springer 2016)

- [2] Michael Kneissl, Tae-Yeon Seong, Jung Han, Hiroshi Amano, "The emergence and prospects of deep ultraviolet light emitting diode technologies", Nature Photonics 13, 233 (2019).
- [3] C.Reich, M.Guttman, M.Feneberg, T.Wernicke, F.Mehnke, C.Kuhn, J.Rass, M. Lapeyrade, S.Einfeldt, A.Knauer, V.Kueller, M.Weyers, R.Goldhahn, M.Kneissl, Appl. Phys. Lett. **107**, 142101 (2015)
- [4] R.G. Banal, M. Funato, Y. Kawakami, Phys. Rev. B **79**, 121308, 2009
- [5] M. Kneissl, T. Kolbe, C. Chua, V. Kueller, N Lobo, J. Stellmach, A. Knauer, H. Rodriguez, S. Einfeldt, Z. Yang, N. M. Johnson, M. Weyers, Semicond. Sci. Technol. **520**, 014036 (2011)
- [6] Y.Taniyasu, M. Kasu, Appl. Phys. Lett. **96**, 221110 (2010)
- [7] H.Y.Ryu, I.G. Choi, J.I. Shim, APEX **6**, 062101 (2013)
- [8] F. Bernardini, V. Fiorentini, D. Vanderbilt, Physical Review B 56 (1997) R10024
- [9] S.Pearton, S. Nakamura, G. Fasol, The Blue Laser Diode: The Complete Story (Springer 1997)
- [10] U.T. Schwarz, M. Kneissl, Phys. Stat. Solidi (RRL) **1**, A44 (2007)
- [11] J.Stellmach, M.Frentrup, F.Mehnke, M. Pristovsek, T.Wernicke, M. Kneissl, J. Cryst. Growth **355**, 59 (2012)
- [12]Y.Zhao, J. Sonoda, C. Pan, S. Brinkley, I. Koslow, K. Fujito, H. Ohta, S. P. DenBaars, S. Nakamura. Appl. Phys. Express **3**, 102101 (2010)
- [13] C.Pan, S. Tanaka, F. Wu, Y.Zhao, J. S. Speck, S. Nakamura, S. P. DenBaars, D. Feezell. Appl. Phys. Express **5**, 062103 (2012)
- [14] K. Balakrishnan, V. Adivarahan, Q. Fareed, M. Lachab, B. Zhang, A. Khan, Jap. J. Appl. Phys. **49**, 040206 (2010).
- [15] M. Lapeyrade, A. Muhin, S. Einfeldt, U. Zeimer, A. Mogilatenko, M. Weyers, and M. Kneissl, Semicond. Sci. Technol. 28, 125015 (2013)

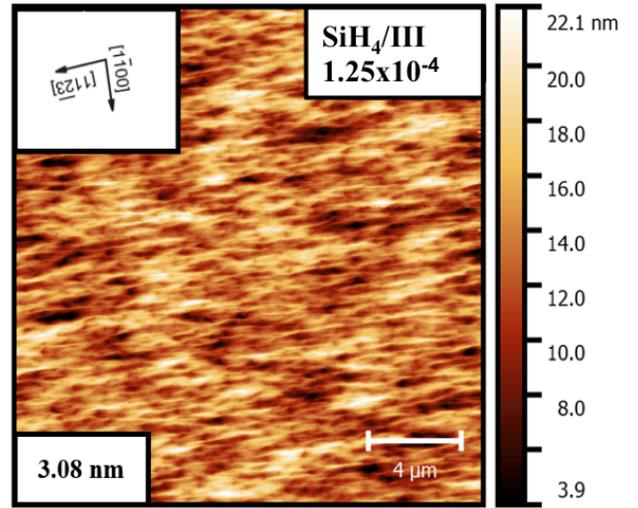
- [16] D.V.Dinh, P.Pampili, P.J.Parbrook *Journal of Crystal Growth* 451, 181-187 (2016)
- [17] Q.Dai, X. Zhang, J.Zhao, H.Luan, Z.Liang, Y. Cui, *Materials Science in Semiconductor Processing* Volume 58, (2017)
- [18] F. Mehnke, L. Sulmoni, M. Guttman, T. Wernicke, and M. Kneissl, *Applied Physics Express*, 012008 (2019)
- [19] H.M Foronda, S. Graupeter, F. Mehnke, J. Enslin, T. Wernicke, M. Kneissl, *Japanese Journal of Applied Physics* 58, SC1026 (2019)
- [20] J.Stellmach, M.Frentrup, F.Mehnke, M. Pristovsek, T.Wernicke, M. Kneissl, *J. Cryst. Growth* **355**, 59 (2012)
- [21] J. Stellmach, F Mehnke, M Frentrup, Ch Reich, J Schlegel, M Pristovsek, T Wernicke, M Kneissl, *Journal of Crystal Growth* **367**, 42 (2013)
- [22] A. Mogilatenko, H. Kirmse, J. Stellmach, M. Frentrup, F. Mehnke, T. Wernicke, M. Kneissl, M. Weyers, *J. Cryst. Growth* **400**, 54 (2014)
- [23] A. Mogilatenko, H. Kirmse, J. Stellmach, M. Frentrup, F. Mehnke, T. Wernicke, M. Kneissl, and M. Weyers, *J. Cryst. Growth* 400, 54 (2014).
- [24] J. Stellmach, F. Mehnke, M. Frentrup, C. Reich, J. Schlegel, M. Pristovsek, T. Wernicke, and M. Kneissl, *J. Cryst. Growth* 367, 42 (2013).
- [25]D. Drouin, A. R. Couture, D. Joly, X. Tastet, V. Aimez, and R. Gauvin, *Scanning*, vol. 29, no. 3, pp. 92–101, May 2007
- [26]G. Kusch, F. Mehnke, J. Enslin, P. R Edwards, T. Wernicke, M. Kneissl, R.W Martin, *Semicond. Sci. Technol.*, vol. 32, no. 3, p. 35020, 2017
- [27]C. Trager-Cowan, A Alasmari, W Avis, Jochen Bruckbauer, PR Edwards, B Hourahine, S Kraeusel, G Kusch, R Johnston, G Naresh-Kumar, RW Martin, M Nouf-Allahiani, E Pascal, L Spasevski, D Thomson, S Vespucci, PJ Parbrook, MD Smith, J Enslin, F Mehnke, M Kneissl, C Kuhn, T Wernicke, S Hagedorn, A Knauer, V Kueller, S Walde, M Weyers, P-M Coulon, PA Shields, Y Zhang, L Jiu, Yipin Gong, RM Smith, T Wang, A Winkelmann, *Photonics Res.*, vol. 7, no. 11, p. B73, 2019
- [28] L. Spasevski, G. Kusch, P. Pampili, V. Z. Zubialevich, D.V. Dinh, J. Bruckbauer, P.R. Edwards, P.J. Parbrook, and R.W. Martin, *J. Phys. D: Appl. Phys.* 54 (2021) 035302
- [29] Luca Sulmoni, Frank Mehnke, Anna Mogilatenko, Martin Guttman, Tim Wernicke, and Michael Kneissl, "Electrical properties and microstructure formation of V/Al-based n-contacts on high Al mole fraction n-AlGaN layers," *Photon. Res.* 8, 1381-1387 (2020)

- [30] F. Mehnke, X. T. Trinh, H. Pingel, T. Wernicke, E. Janzen, N. T. Son, and M. Kneissl, *J. Appl. Phys.* 120, 145702 (2016).
- [31] D.A. Browne, E.C. Young, J.R. Lang, C.A. Hurni, J.S. Speck, *Journal of Vacuum Science & Technology A* 30, 041513 (2012)
- [32] D.C. Look, *Electrical Characterization of GaAs Materials and Devices*, John Wiley and Sons Ltd. (1989)
- [33] I. Bryan, Z. Bryan, S. Washiyama, P. Reddy, B. Gaddy, B. Sarkar, M.H. Breckenridge, Q. Guo, M. Bobea, J. Tweedie, S. Mita, D. Irving, R. Collazo, Z. Sitar *APPLIED PHYSICS LETTERS* 112, 062102 (2018)
- [34] F. Mehnke, T. Wernicke, H. Pingel, C. Kuhn, C. Reich, V. Kueller, A. Knauer, M. Lapeyrade, M. Weyers, and M. Kneissl, *Appl. Phys. Lett.* 103, 212109 (2013).
- [35] J.L. Farvacque, *PHYSICAL REVIEW B*, 62, 4 (2000)
- [36] W.T. Read Jr. LXXXVII. Theory of dislocations in germanium, *Philosophical Magazine Series 7*, 45:367, 775-796, (1954)
- [37] J.L. Farvacque, Z. Bougrioua, and I. Moerman, *PHYSICAL REVIEW B*, 63, 115202, (2001)
- [38] S. F. Chichibu, H. Miyake, Y. Ishikawa, M. Tashiro, T. Ohtomo, K. Furusawa, K. Hazu, K. Hiramatsu, and A. Uedono, *J. Appl. Phys.* 113, 213506 (2013).
- [39] K.B. Nam, J. Li, M.L. Nakarmi, J.Y. Lin, H.X. Jiang, *Appl. Phys. Lett.* 81 (2002) 1038.
- [40] K.X. Chen, Q. Dai, W. Lee, J.K. Kim, E.F. Schubert, W. Liu, S. Wu, X. Li, J.A. Smart, *Appl. Phys. Lett.* 91 (2007) 121110.
- [41] A. Uedono, K. Tenjinbayashi, T. Tsutsui, Y. Shimahara, H. Miyake, K. Hiramatsu, N. Oshima, R. Suzuki, S. Ishibashi, *J. Appl. Phys.* 111 (2012) 013512
- [42] A. Kakanakova-Georgieva, D. Nilsson, X. T. Trinh, U. Forsberg, N. T. Son, and E. Janzen *Appl. Phys. Lett.* 102, 132113 (2013);
- [43] Joshua S. Harris, Benjamin E. Gaddy, Ramón Collazo, Zlatko Sitar, and Douglas L. Irving *Phys. Rev. Materials* 3, 054604 (2019)

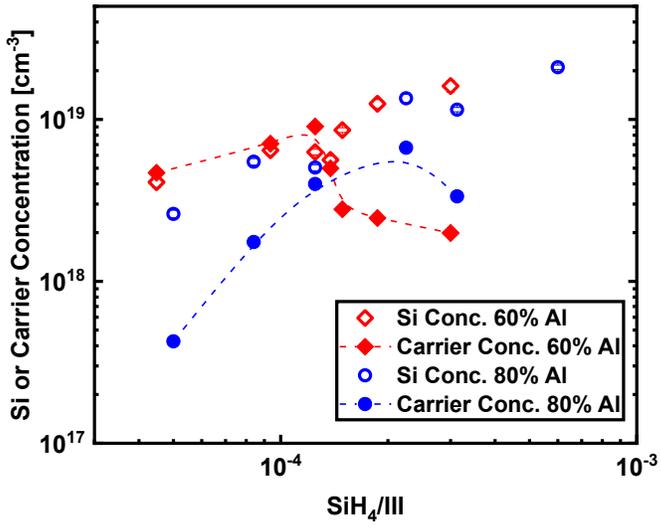
(a)



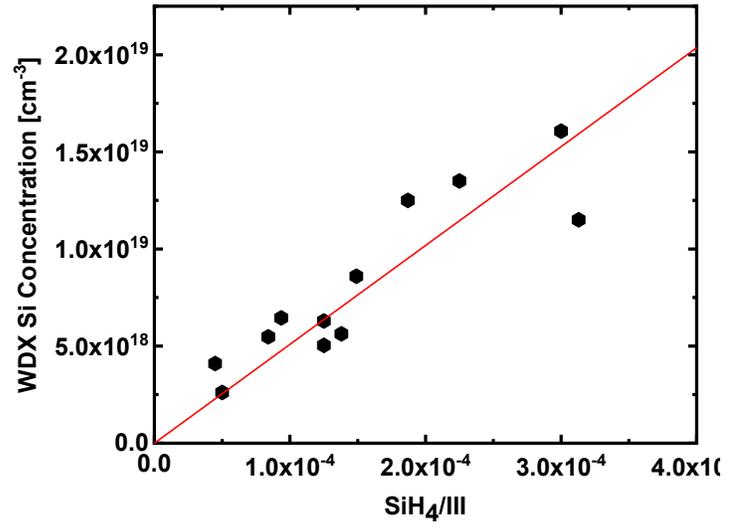
(b)



(a)



(b)

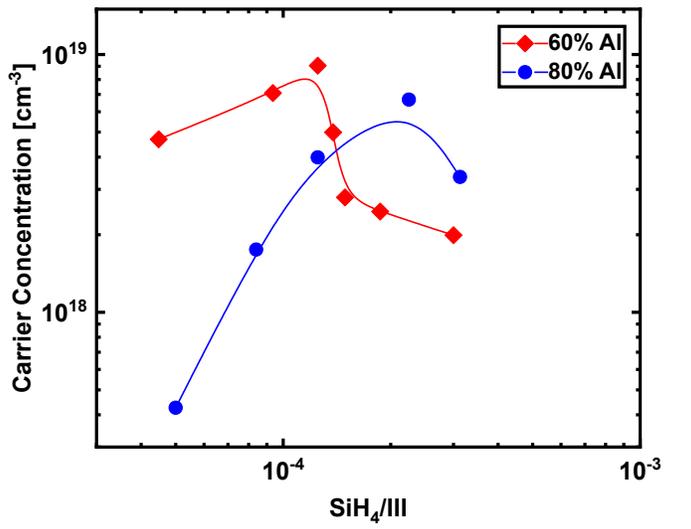


(a)

Si concentration from WDX fit [cm^{-3}]

5.1×10^{18}

5.1×10^{19}



(b)

Si concentration from WDX fit [cm^{-3}]

5.1×10^{18}

5.1×10^{19}

