

# Human Red Blood Cell Behavior under Homogeneous Extensional Flow in a Hyperbolic-Shaped Microchannel

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## ABSTRACT

*It is well known that certain pathological conditions result in a decrease of red blood cells (RBCs) deformability and subsequently can significantly alter the blood flow in microcirculation, which may block capillaries and cause ischemia in the tissues. Microfluidic systems able to obtain reliable quantitative measurements of RBC deformability hold the key to understand and diagnose RBC related diseases. In this work, a microfluidic system composed of a microchannel with a hyperbolic-shaped contraction followed by a sudden expansion is presented. We provide a detailed quantitative description of the degree of deformation of human RBCs under a controlled homogeneous extensional flow field. We measured the deformation index (DI) as well as the velocity of the RBCs travelling along the centerline of the channel for four different flow rates and analyze the impact of the particle Reynolds number. The results show that human RBC deformation tends to reach a plateau value in the region of constant extensional rate, the value of which depends on the extension rate. Additionally, we observe that the presence of a sudden expansion downstream of the hyperbolic contraction modifies the spatial distribution of cells and substantially increases the cell free layer (CFL) downstream the expansion plane similarly to what is seen in other expansion flows. Beyond a certain value of flow rate, there is only a weak effect of inlet flow rates on the enhancement of the downstream CFL. These in vitro experiments show the potential of using microfluidic systems with hyperbolic-shaped microchannels for both the separation of the RBCs from plasma and to assess changes in RBC deformability in physiological and pathological situations for clinical purposes. However, the selection of the geometry and the identification of the most suitable region to evaluate the changes on the RBC deformability under extensional flows are crucial if to be used as an in vitro clinical methodology to detect circulatory diseases.*

**Keywords:** Extensional flows, red blood cells, hyperbolic microchannel, deformation index.

## 1 INTRODUCTION

Erythrocyte deformability refers to the ability of the red blood cells (RBCs) to deform when subjected to certain flow conditions in both *in vivo* and *in vitro* environments. The major determinants of the deformability include geometry and size of the microvessels, mechanical properties of the cell membrane and its cytoskeleton and intracellular viscosity. For example, RBCs elongate significantly when they pass through capillaries, the diameter of which is smaller than the diameter of the RBCs at rest [1-5]. The deformability (usually reported as a deformation index, DI) is an important property for the delivery of oxygen to the body and a decrease in RBC deformability can have serious consequences leading to health problems [5]. There are several studies examining the behaviour of RBCs in response to a deforming force (or stress) and demonstrating the clinical relevance of this phenomenon in several blood cells related diseases, such as malaria [4, 6], and diabetes mellitus (DM) [5, 7, 8] in which the deformability is significantly reduced. The capability of the cells to deform has, therefore, the potential to be used as a

biophysical marker, without requiring the costly labeling and sample preparation typical of most common biochemical markers [9].

During the last decade, the deformation capability of RBCs and its associate effects such as Fahraeus effect, margination and plasma skimming have been applied to develop microfluidic systems for blood cells separation and analysis [10]. Faivre et al. [11] showed that varying geometrical parameters such as the width and length of a constriction can promote a dramatic enhancement of the cell-free layer (CFL) downstream of the constriction region. More recently, Leble et al. [12] have also demonstrated the formation of a CFL around the apex of a confluence and its propagation in a straight line downstream of this region. Additionally, Fujiwara et al. [13] have determined the motion of both normal and hardened RBCs in a microchannel with a stenosis using a confocal system. This study has shown evidence that not only it is possible to create an artificial CFL under appropriate hemodynamic and geometrical conditions but also the CFL thickness is strongly influenced by the RBC deformability. These findings have been promoting the design of new alternative biomedical microdevices able to separate blood cells from plasma and also to detect cellular-scale changes, which can be insightful for medical diagnosis. In particular Yang et al. [14] and Shevkoplyas et al. [15] used biomimetic approaches to separate blood cells from plasma, and leukocytes from whole blood, respectively. Additionally, Shevkoplyas et al. [16] also proposed a microfluidic device for assessing RBC deformability in a microchannel network with a topology similar to that of the real microcirculation and showed the ability to detect small changes of RBC deformability on the blood flow in a microchannel network may prove a viable clinical instrument in early detection of blood diseases. More recently, Hou et al. [6] have proposed a biomimetic separation device to separate normal and malaria infected RBCs.

Investigations on human RBC deformability have been performed using a variety of techniques, including optical tweezers, micropipette aspiration, rheoscope and ektacytometer [17, 18]. These and other techniques were thoroughly reviewed by Kim et al. [19] and are useful in measuring the RBC deformation under shear flow or the rheological properties of the RBC membrane, but are generally labor intensive, time consuming and involve a typically difficult process of manipulation. Among these methods, the micropipette aspiration is the most popular single-cell technique to measure the mechanical properties of the RBC membrane [18]. Briefly, the RBC membrane is aspirated by a negative pressure through a glass micropipette, having inner diameter much smaller than the RBC average size, forcing the micropipetted RBCs to adopt an extended “tongue” shape where its length is measured under a microscope. A major drawback of this technique is that direct contact between the cell surface and the wall cannot be avoided and the extensional rate is not constant in space and time. Hence, the mechanical properties of the RBCs measured with this method are affected, for example, by the pipette shape, pressure difference and viscosity ratio.

The methodology adopted in this paper, using a microfluidic device with a hyperbolic-shaped contraction, is able to assess RBC deformability under extensional flow in a region of homogeneous extension rate. The recent progresses in microfabrication, optics and image analysis techniques made it possible to develop microfluidic devices able to precisely measure and control the dynamical deformation of cells flowing through microchannels. The distinct advantage of the microfluidic devices to test a large number of cells using only a small amount of blood has prompted a vast amount of research in this field [8, 10, 13, 16, 20-22]. However, most of the proposed devices focus mainly on the effect of shear flow alone. Some examples are the measurement of the RBC cellular trajectories and deformation under a transient high shear stress in microchannels with a rectangular cross-section [21, 22] and the determination of cell deformability of a single RBC flowing in a microfluidic device with a microchannel thickness comparable to the RBC size [20]. Studies focusing on the effect of extensional flow field on the RBC flow and deformability are scarce despite extensional flow often being present in the microcirculation, particularly when there is a sudden change in geometry, e.g. in stenosis and in microvascular networks composed of short irregular vessel segments linked by numerous bifurcations. In

converging flows such as those found in microstenoses, extensional flow can be generated due to the dramatic change of velocity when the fluid flows from a wide to a narrow region of the channel. The behavior of an RBC under extensional flow is known to be different from that under shear flow. For instance, in the case of extensional flow, an RBC does not show tumbling motion as in shear flow [23]. However, details of these differences have not been clarified fully in former studies.

In the present work, we use a microfluidic hyperbolic-shaped contraction to achieve an extensional flow, in which RBCs are submitted to high strains and a range of strain-rates. The rationale behind the use of hyperbolic contractions is their ability to impose a nominally constant strain rate ( $\dot{\epsilon} = du_x / dx = \text{constant} \neq 0$ ) along the centerline as discussed by James [24] for an axisymmetric configuration. Relevant works in the context of blood flow are those of Sousa et al. [25] and Lee et al. [26] who studied the effect of the extensional flow in hyperbolic converging microchannels using artificial polymer-based blood analog fluids and *in vitro* rabbit blood, respectively. Sousa et al. [25] used hyperbolic contractions of different Hencky strains to show that two well-established blood analogue fluids with identical steady-shear viscosity, but different extensional rheology, behave very differently when submitted to conditions equivalent to the microcirculation. The blood analog fluids considered in that study were polymeric solutions and therefore no study on the effect of RBC deformability under extensional flow was undertaken. Lee et al. [26] has shown, using rabbit RBCs, that the extensional flow was more efficient in inducing deformation than shear flow. Moreover, Lee et al. [26] have only performed RBC deformability measurements in one single region around the middle of the hyperbolic contraction and assumed that the nominal strain rate was achieved in this region. However, Oliveira et al [27] have shown that in planar microchannels the strain rate along the contraction is not truly homogeneous. The authors examined the flow kinematics in detail using micro-particle image velocimetry ( $\mu\text{PIV}$ ) and showed that the velocity profile in the contraction region deviates from linear due to wall, entrance and end effects resulting from the low aspect ratios typical of microfluidic devices and from the sharp transitions to and from the contraction region. Moreover, the cell flow behavior may be affected by variables such as cell shape, particle Reynolds number (ratio of inertial to viscous forces), Peclet number (ratio of advection rate and diffusion rate) and Stokes number (ratio of the relaxation time of the cell to the characteristic time of the flow), and depending on these conditions cells may no longer follow the fluid streamlines [28]. Therefore, to be able to use this type of experiment as a biophysical marker in a diagnostic context it is crucial that the actual flow field is well understood. In the present paper we examine the evolution of RBC deformability, quantified in terms of DI values, along the full length of the hyperbolic converging microchannel, including upstream and a sudden expansion release plane downstream, and simultaneously evaluate the cell velocity, quantify the extension rate and identify the region of homogeneous strain-rate. This study is important not only to understand the RBC motion and mechanical properties under the effect of extensional flows but also to help clinicians decide in which region to perform analysis of RBC deformability under pathological conditions, for which reliable quantitative techniques are still lacking.

The main objective of this study is to investigate how human RBCs respond under a controlled extensional flow field. For this purpose RBC deformation index (DI) was measured along the whole microchannel, with a hyperbolic-shaped contraction followed by a sudden expansion, for a range of imposed volumetric flow rates. By using a high-speed video microscopy system combined with image analysis technique, we were able to provide detailed quantitative description of the RBC DI for a range of imposed volumetric flow rates. Additionally, cell tracking was performed to quantify the velocity of the cells along the channel and evaluate the geometrical effect of the hyperbolic contraction-sudden expansion on the spatial distribution of RBCs and the size of the cell-free layer (CFL). Finally, we also show the benefits and the potential applications of this unique approach to assess RBC deformability and

how it can be integrated into a diagnostic microfluidic system to simultaneously deform and separate cells.

## 2 MATERIALS AND METHODS

### 2.1 Working Fluids and Microchannel Geometry

The working fluid used in the present study was composed of Dextran 40 (Dx40) containing ~1% of human RBCs (i.e., hematocrit, Hct = 1%). The density and viscosity of Dextran 40 are  $1046 \text{ kg m}^{-3}$  and  $4.5 \times 10^{-3} \text{ Pa s}$ , respectively. Blood was collected from a healthy adult volunteer, and EDTA (ethylenediaminetetraacetic acid) was added to the collected samples to prevent coagulation. The blood samples were submitted to washing and centrifuging processes to separate the RBCs from the plasma and buffy coat. The RBCs were then added to Dx40 to make up a suspension with the required RBC concentration, and were subsequently stored hermetically at  $4^\circ\text{C}$  until the experiments were performed. All procedures in this work were carried out in compliance with the Ethics Committee on Clinical Investigation of Tohoku University.

The microchannels containing the hyperbolic contraction were produced in polydimethylsiloxane (PDMS) using standard soft-lithography techniques from a SU-8 photoresist mold. The molds were prepared in a clean room facility by photo-lithography using a high-resolution chrome mask. The geometry and real dimensions of the micro-fabricated channels are shown in Figure 1.

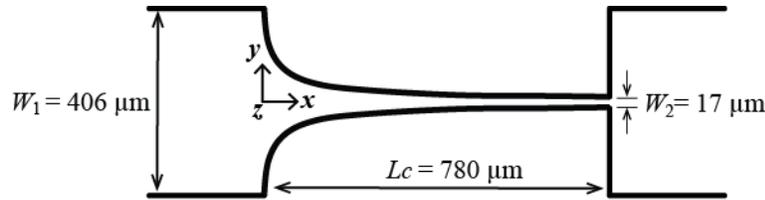


Figure 1. Schematics of the hyperbolic contraction geometry used in the experiments. The channel depth,  $d$ , was  $60 \mu\text{m}$  and the width of the upstream and downstream channels were the same,  $W_1 = 406 \mu\text{m}$ . The minimum width in the contraction region was  $W_2 = 17 \mu\text{m}$ , defining a total Hencky strain of  $\epsilon_H = \ln(W_1/W_2) = 23.9$ .

For the microfluidic experiments, the channels were placed on the stage of an inverted microscope (IX71, Olympus, Japan) with an objective lens with a magnification of  $20\times$  (Optical zoom  $1.6\times$ ). The temperature of the stage was adjusted by means of a thermo plate controller (Tokai Hit, Japan) to  $37^\circ\text{C}$ . The flow rate of the working fluids was controlled using a syringe pump (KD Scientific Inc., USA), and four different flow rates were examined ranging from  $0.11 \text{ ml/h}$  to  $4.20 \text{ ml/h}$ , corresponding to a maximum Reynolds number ( $Re$ ) based on the upstream variables of 1.8. The images of the flowing RBCs were captured using a high speed camera (Phantom v7.1, Vision Research, USA) for three regions: the beginning of the contraction region, the end of the contraction region and the expansion region. The exposure time was  $2 \mu\text{s}$  and the videos were captured at frame rates ranging from 4800 frames/s to 13000 frames/s, depending of the region of the channel and the flow rate under study. Finally, these images were transferred to a computer to be analyzed. An illustration of the experimental set-up is shown in Figure 2 and relevant experimental parameters are summarized in Table 1.

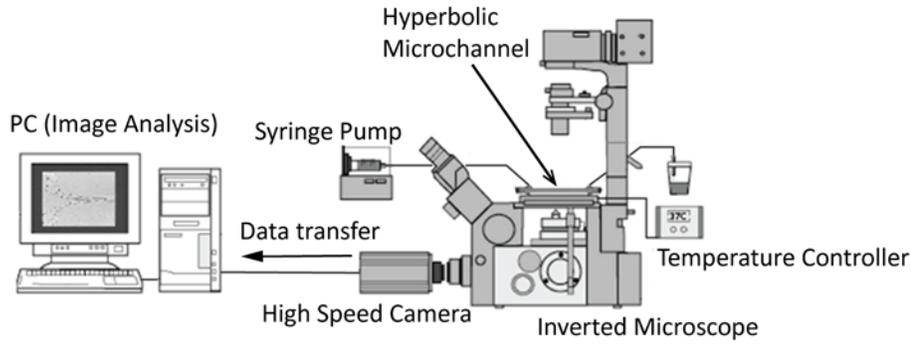


Figure 2: Experimental set-up consisting of an inverted microscope (IX71, Olympus, Japan), a high speed camera (Phantom v7.1, Vision Research, USA), a syringe pump (KD Scientific Inc., USA) and a thermo plate controller (Tokai Hit, Japan).

Table 1. Summary of experimental parameters

Experimental parameters	
Width of microchannel upstream and downstream	406 $\mu\text{m}$
Minimum width of contraction	17 $\mu\text{m}$
Length of contraction region	780 $\mu\text{m}$
Depth of microchannel	60 $\mu\text{m}$
Inlet flow rate	0.11, 1.11, 2.27 and 4.2 ml/h
Mean velocity of the fluid	1.56, 1.56, 31.5 and 58.2 mm/s
Mean diameter of human RBC at rest	8 $\mu\text{m}$
Shear viscosity of Dextran 40	$4.5 \times 10^{-3}$ Pa s
Density of Dextran 40	1046 $\text{kg}/\text{m}^3$
Hematocrit of working fluid	1 %
Temperature	37 $^{\circ}\text{C}$
Magnification (M)	20 $\times$ (Optical zoom 1.6 $\times$ )
Frame rate of camera	4800 - 13000 frames/s
Exposure time	2 $\mu\text{s}$
Frame interval	77 – 208 $\mu\text{s}$

## 2.2 Image Analysis

Image analysis was used to determine the cells trajectories and evaluate their velocity, deformability, and finally the thickness of the cell free layer upstream and downstream of the contraction region. The first step of this process involves converting the captured video files to a sequence of static images (known as a stack), with a resolution of  $800 \times 600$  pixels at frame intervals of 208  $\mu\text{s}$  for the beginning of the contraction part images and expansion plane images, and a resolution of  $800 \times 200$  pixels at frame intervals of 77  $\mu\text{s}$  for the ending part of the contraction region images. A typical image obtained is shown in Figure 3a. Then, in order to reduce the dust and static artifacts in the images, a background image (Figure 3b) was created from the original stack images by averaging each pixel over the sequence of static images using an image handling software ImageJ (NIH) function called *Z project*, and then subtracted

from the stack images [29]. This process eliminates all the static objects from the original images including the microchannel walls, which resulted in images having only the flowing RBCs visible (Figure 3c). To enhance the image quality, image filtering such as *Median* operation with a mask size of 3×3 pixels, and *Brightness/Contrast* adjustment was also applied. Finally, the greyscale images were converted to binary images adjusting the threshold level (Figure 3d). In this study, *Otsu* threshold method was applied first and when required the level was refined manually for optimal binarization. In general, the pixels with intensity levels in the range of 0-70 were set to be 0 (black) and the pixels with intensity levels greater than 70 were set to be 255 (white); although the level was slightly adjusted depending on the quality of each image stack, the differences are minimal. This segmentation process yields regions of interest (RBCs) as black ellipsoidal objects against a white background. Next, the flowing cells in the binary images were measured frame by frame automatically by *Analyze Particles* function in ImageJ. Various measurement parameters, such as *Area* and *Circularity* of the cells, were appropriately pre-set to ignore the apparent deviant objects (eg. out-of-focus cells, aggregated cells, etc.). The major output results of this measurement are the major and minor axis lengths of the RBCs and the *x-y* coordinates of their centroid.

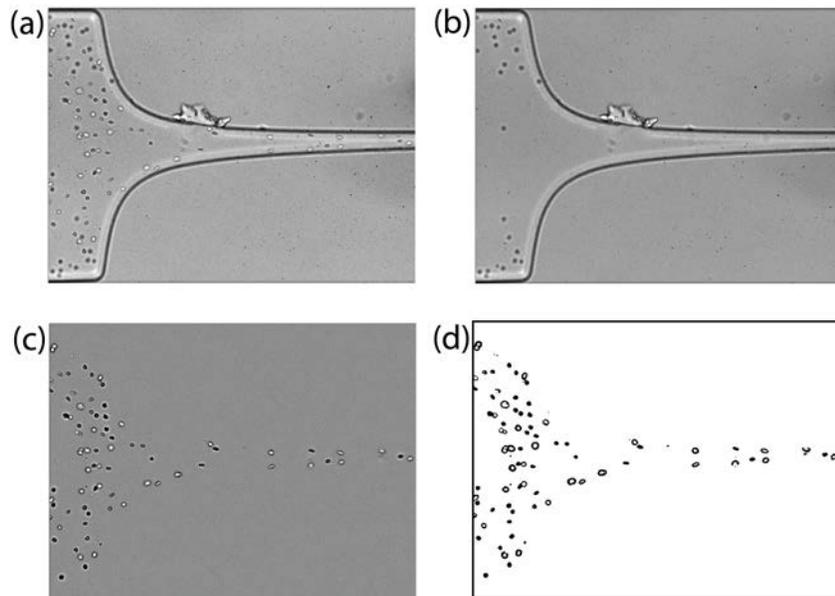
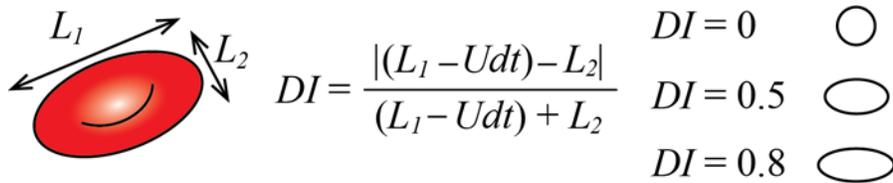


Figure 3. Images analysis sequence: a) original image in which moving cells as well as microchannel boundaries and other static objects are visible, b) background image containing only static objects, c) original image after background subtraction showing only moving objects, and d) final binary image.

Using this set of data, the deformation index (DI) of all the measured cells was calculated and stored together with the cells positions given by their *x-y* coordinates. In this study, DI is defined by the equation shown in Figure 4, where  $L_1$  and  $L_2$  refer to the major (primary) and minor (secondary) axis lengths of the ellipse best fitted to the cell. The correction  $Udt$  term refers to the traveling length of RBC that is calculated by the RBC velocity times the exposure time of high speed camera.



$$DI = \frac{|(L_1 - Udt) - L_2|}{(L_1 - Udt) + L_2}$$

$DI = 0$      ○  
 $DI = 0.5$     ○  
 $DI = 0.8$     ○

Figure 4: Definition of the deformation index,  $DI = \frac{|(L_1 - Udt) - L_2|}{(L_1 - Udt) + L_2}$ , where  $L_1$  and  $L_2$  are the major (primary) and minor (secondary) axis lengths of the ellipse best fitted to the cell,  $Udt$  is the traveling length of RBC that is calculated by the RBC velocity ( $U$ ) times the exposure times of high speed camera ( $dt$ ).

For each flow rate, the trajectories of about 30 individual RBCs were tracked manually using ImageJ plug-in, MTrackJ, in order to examine the actual velocity of the flowing cells [30-32]. Sample cells were selected at various positions in the y-axis in the upstream region of the microchannel to obtain a representative illustration of the trajectories as shown in Figure 5. We should note that the same cell cannot be tracked all the way from upstream to downstream of the contraction region, and Figure 5 is a combination of three separate sets of images corresponding to different axial regions. However, careful superposition and selection of the cells to match the same position in consecutive images allow us to show nearly full cell path lines.

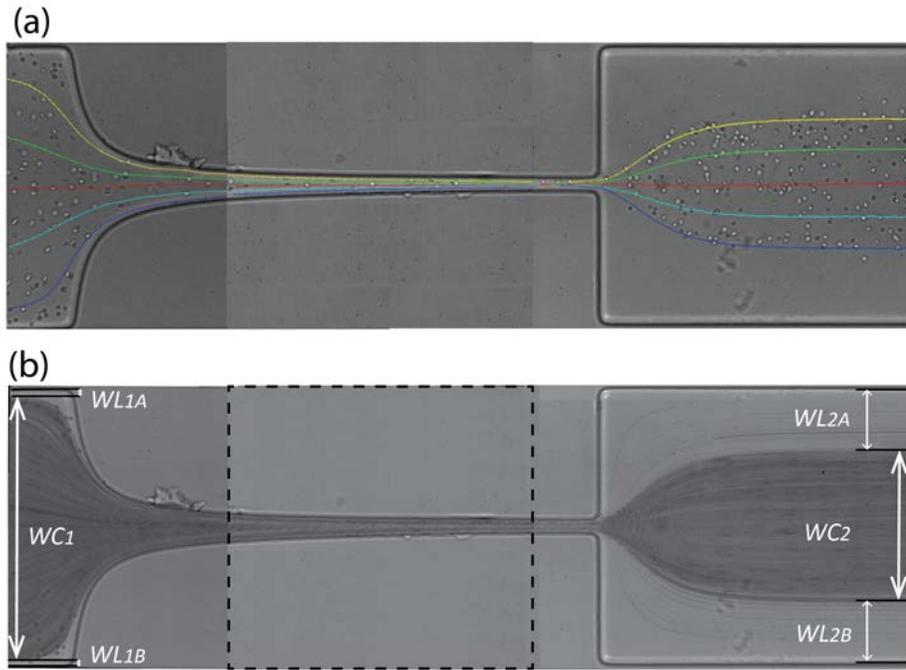


Figure 5. a) RBC trajectories manually tracked by MTrackJ plug-in from ImageJ. b) Image obtained after post-processing using the minimum intensity level option in ImageJ, *Zproject* operation for the measurement of the cell free layer and relevant variables:  $WC_1$  and  $WC_2$  are the widths of RBC core upstream and downstream of the hyperbolic region, respectively.  $WL_{1A}$  and  $WL_{1B}$  are the widths of CFL upstream of the contraction region, and  $WL_{2A}$  and  $WL_{2B}$  are the widths of CFL downstream of the contraction region. Each image is a combination of three separate images covering different parts of the channel. The dashed lines indicate the boundaries between images.

The cells trajectories clearly highlight the enhancement of the cell free layer downstream of the expansion on the flow. To quantify this enhancement, we have measured the width of the CFLs in the upstream and downstream regions of the microchannel in the following way. First, in order to identify the boundary between RBC core and CFLs, minimum intensity level option in ImageJ *Zproject* function was applied to the original images (stack). This operation generates a new image consisting of the minimum intensity level for each pixel, which distinctively reveals the boundary between the RBC core and the CFLs as shown in Figure 5b. The width of the CFLs and core regions can then be easily evaluated. In our study, the same process was applied for four different flow rates (0.11ml/h, 1.11ml/h, 2.27ml/h and 4.2ml/h) and each CFL and core widths were manually measured using ImageJ. The relevant variables are the widths of RBC core upstream and downstream of the hyperbolic region ( $WC_1$  and  $WC_2$  respectively), the widths of CFL upstream of the contraction region ( $WL_{1A}$  and  $WL_{1B}$ ) and the widths of CFL downstream of the contraction region ( $WL_{2A}$  and  $WL_{2B}$ ) (cf. Figure 5b).

### 3 RESULTS AND DISCUSSION

#### 3.1 Cell velocity and strain rate

The cells traveling along the centerline of the microchannel were manually tracked and their local velocity was determined for four different volumetric flow rates. Figure 6 shows the averaged cell velocity for each 50 $\mu\text{m}$  axial region throughout the central region of the microchannel for the four flow rates. The overall trend for all flow rates is similar. As the cells approach the contraction ( $x/L_c = 0$ ), a cell acceleration is observed due to entrance effects. Within the contraction region, as the width of the channel gets smaller, the cell velocity increases significantly, attaining a maximum value just before the expansion plane ( $x/L_c = 1$ ). The maximum velocity measured is, in most cases, slightly below the actual fully developed value expected for single phase flows. It is clear from the velocity profiles in Figure 6 that the strain rate is not constant for the whole length of the contraction. However, there is a region where the velocity increases almost linearly with the axial position and consequently we can define a constant extension rate in this region. For this region, we can thus define a single characteristic strain rate ( $\dot{\epsilon}$ ) that can be calculated as the slope  $du_x/dx$ , where  $u_x$  refers to the cell axial velocity. The values of strain rate obtained in this way are comparable (within 10%) to the nominal strain rate values calculated for single fluid flow by  $\dot{\epsilon}_{nom} = \frac{U_c - U_1}{L} = Q \frac{w_1 - w_c}{Lh w_1 w_c}$ . However, it should be noted that when high particle

Reynolds numbers ( $Re_p$ ) are attained, differences relative to the expected fluid-based nominal strain rate value may be enhanced and should be accounted for. As the cells exit the contraction region, they decelerate and the velocity eventually approaches a constant value, which is consistent with fully-developed flow for a rectangular channel with this aspect ratio.

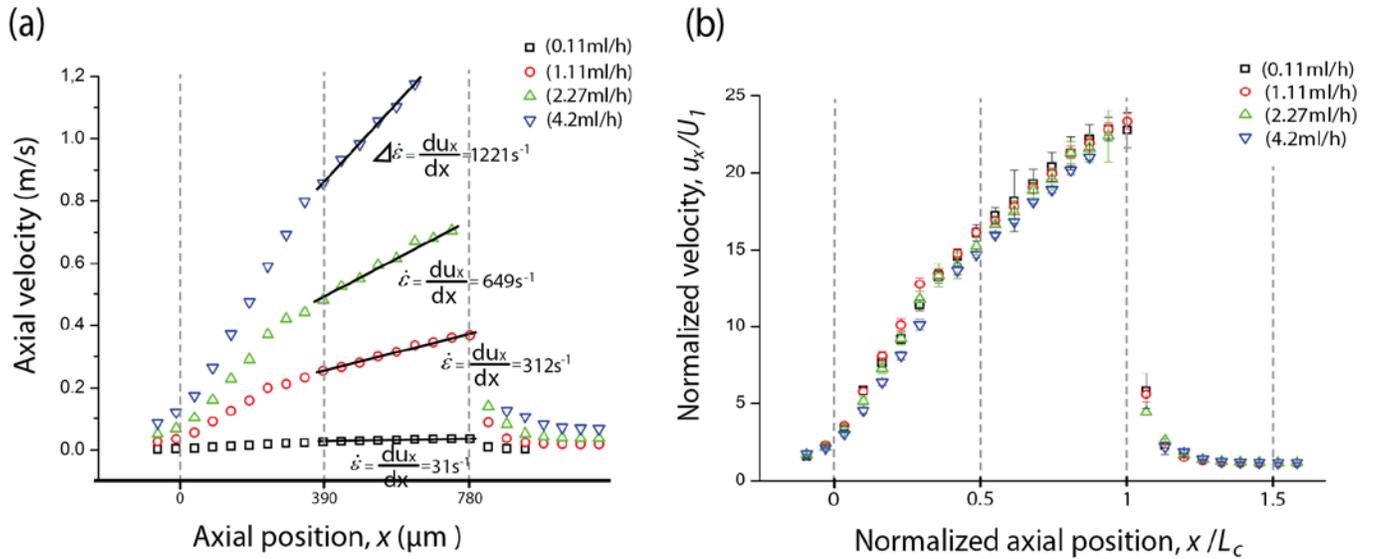


Figure 6. a) Axial velocity profiles of flowing RBCs along the centerline for four flow rates ( $Q=0.11\text{ml/h}$ ,  $Q=1.11\text{ml/h}$ ,  $Q=2.27\text{ml/h}$  and  $Q=4.2\text{ml/h}$ ). The solid lines represent the fit used to determine the strain rate in the linear region, calculated from the slope  $du_x/dx$ . b) Normalized axial velocity profile  $u_x/U_1$ , where  $u_x$  refers to an instantaneous cell velocity average, and  $U_1$  is the cross-section overall mean velocity based on the large channel,  $U_1 = Q/(d W_1)$ . The dashed lines indicate the start ( $x/L_c = 0$ ), mid ( $x/L_c = 0.5$ ) and end ( $x/L_c = 1$ ) of the contraction region.

### 3.2 Deformation Index

In this section, we focus on the deformation of the RBCs travelling near the microchannel centerline ( $y = 0$ ). As illustrated in Figure 7(a), 12 regions were pre-defined in order to analyze the degree of DI along the centerline. Each rectangular region is  $50\mu\text{m}$  long and  $20\mu\text{m}$  high and all the cells captured in each region were measured and the averaged DI was calculated as described in section 2.2.

Far upstream of the contraction region, the orientation of the RBCs is affected mainly by shear stress induced by the out-of-plane walls (top and bottom walls on the  $z$ -direction). According to Omori et al. [33, 34], the RBC orientation tends to converge to a certain direction when the cells are submitted to strong shear for a sufficiently long-time. Therefore, in the present case, most of the RBCs tend to direct parallel to the out-of-plane walls, showing a nearly circular shape, which is desirable for two-dimensional measurements of RBC DI, as the errors caused by the RBC orientation are minimized. Moreover, the shear caused by these nearby walls may induce the RBC to show tank-treading motion with its flattened plane parallel to the out-of-plane walls. To examine whether or not the RBCs are actually in the tank-treading regime, it is relevant to know if the RBCs are travelling steadily with a similar orientation. There are two important dimensionless numbers that need to be considered in this context, the viscosity ratio of

inside to outside fluid  $\left(\lambda = \frac{\mu_{\text{cell}}}{\mu}\right)$  and the capillary number,  $Ca$ , defined as the ratio of the fluid viscous forces causing the deformation and the membrane elastic forces that resist deformation:

$$Ca = \frac{\mu a \dot{\gamma}}{G_s} \quad (1)$$

where  $\mu$  is the shear viscosity of Dextran 40 solution,  $\dot{\gamma}$  is the shear rate,  $a$  is the characteristic length of the undeformed RBC and  $G_s$  is the shear elastic modulus [33, 34]. The shear rate was estimated as  $2u_x/H$ , where  $u_x$  is the average axial velocity at the centerline of the defined region and  $H$  is  $\min(W, h)$  where  $W$  is the width of the microchannel and  $h$  is the distance between upper and lower walls. The  $Ca$  values were calculated for several regions along the centerline of the microchannel as a function of flow rate and are presented in Table 2 assuming  $a=2.82\mu\text{m}$  [35] and  $G_s=4\mu\text{N/m}$  [36]. It is generally agreed that in shear flow, suspended RBCs show smooth tank-treading motion when  $Ca$  is larger than about 1 [33, 34] and the suspending medium viscosity is sufficiently high, with the viscosity ratio of inside to outside fluid being about unity [37]. In the present study, we use Dextran 40, which yields a viscosity ratio about 1.0-1.5. However, it is not been clarified in previous studies how the above criterion for tank-treading motion changes with flow type (recall that the flow at the centerline of the contraction is extensionally-, rather than shear-dominated). Nevertheless, we believe the instantaneous Capillary number (cf. Table.2), is still insightful for a broad understanding of the RBCs motion. Except for the lowest flow rate ( $Q = 1.11\text{ml/h}$ ), the results obtained from these examined regions show sufficiently high  $Ca$  values, in particular for  $x/L_c \geq 0.5$ .

Table 2. Capillary number,  $Ca$ , for three regions at the microchannel for different flow rates.

Regions*	Flow rates			
	$Q = 0.11\text{ml/h}$	$Q = 1.11\text{ml/h}$	$Q = 2.27\text{ml/h}$	$Q = 4.2\text{ml/h}$
A	0.29	3.12	6.14	9.94
C	3.13	33.94	59.34	98.02
F	11.00	110.08	211.63	383.67

\*Note that the regions correspond to those depicted in Fig.7.

Figure 7(b) presents the average DI for the 12 regions and for four different flow rates. The average cell number analyzed for each region is 431. The error bar means 95 % confidence interval. As can be seen, far upstream of the contraction region, the cells at the centerline are nearly circular, exhibiting a DI close to zero. Note that, for rigid cells and for cells at rest, the cell shape should be nearly circular corresponding to  $DI = 0$  [38]. Under flow, the cells change from a circular to an elliptical shape and they become increasingly elongated as they flow through the contraction. Generally, after the mid-point of the contraction part ( $x/L_c \geq 0.5$  corresponding to Region D in Figure 7a), DI tends to reach a maximum plateau value, with the DI values staying fairly constant between regions E and H. Hence, these results show that, in the region just after the contraction mid-point, human RBCs deformability can reach a plateau value in a constant extensional rate, indicating that the RBCs travelling along the centerline do not change shape. It is also noted that the DI increases substantially with the flow rate as a consequence of the higher strain rate to which the RBCs are submitted. At Region I, DI decreases dramatically and the RBCs start regaining their near circular shape, and when the cells reach Region K, corresponding to about 256  $\mu\text{m}$  downstream of the expansion plane, DI remains approximately constant. The data presented here suggests that the representative DI values depend on the region of the hyperbolic microchannel where the measurements are undertaken and that these have to be carefully selected for meaningful and reliable analysis.

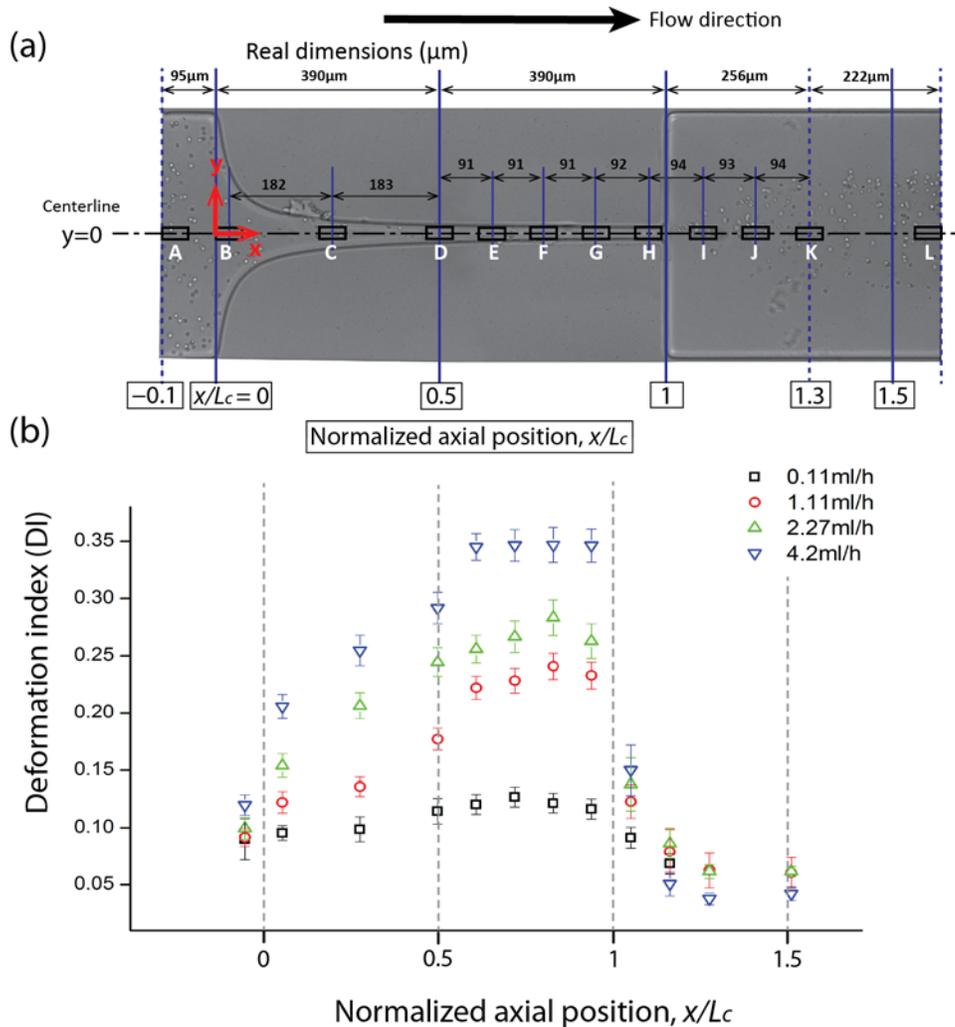


Figure 7. Evaluation of the deformation of the flowing RBCs at the centerline of the hyperbolic microchannel. (a) Illustration of the 12 regions used to analyze the RBCs deformation index. (b) Average DI in the 12 regions at four different volumetric flow rates:  $Q = 0.11\text{ml/h}$ ,  $Q = 1.11\text{ml/h}$ ,  $Q = 2.27\text{ml/h}$  and  $Q = 4.2\text{ml/h}$ . The measured values are expressed as the means  $\pm$  standard deviation according to a t-test analysis at a 95% confidence interval.

It is clear in Figure 7b, that under strong extensional flow, DI is strongly flow rate dependent, with the deformation levels increasing monotonically with flow rate in the range studied. The results obtained under shear flows, by Zhao et al. [21], have shown that under several flow rates RBCs elongation reached the same maximum value and could not deform any further. Hence, under shear flow and beyond a certain value, DI is flow rate independent, in contrast with the results presented in Figure 7b. Based on the analysis of the effect of the particle Reynolds number shown below, we believe a similar behavior occurs in extensional flow, but due to technical limitations of our high-speed micro-visualization system we were only able to obtain good quality images up to a flow rate of  $Q = 4.2$  ml/h. In a further study, using an upgraded micro-visualization system, it would be interesting to elucidate if there is indeed a limit on RBC maximum deformability under extensional flow and if cell lysis is likely to occur. Moreover, a previous study performed by Lee et al. [19], by using rabbit RBCs, have investigated the RBC DI in a single position at the entrance region of a hyperbolic contraction. By using an automated rheoscope, Dobbe et al. [39] have shown that the deformability indexes of rabbit RBCs are higher than human RBCs and consequently the degree of deformation of human RBCs in hyperbolic microchannels may differ from rabbit RBCs. Hence, in contrast with the study performed by Lee et al. [19], here we provide a detailed quantitative description of human RBC DI along the whole hyperbolic microchannel and we were able to observe that the human RBCs tend to reach a converged DI value in the region just after the contraction mid-point.

In order to get further insight into the RBC behavior under flow at the centerline of the hyperbolic microchannel, DI is plotted as a function of the particle Reynolds number (see Fig 8). The DI values are those used for Figure 7b, that is, the average DIs of 12 regions at four different flow rates. The particle Reynolds number ( $Re_p$ ) is useful to assess the behavior of flowing cells suspended in the working fluid taking into account the inertial and viscous forces occurring by interactions of the cells and the fluid.  $Re_p$  can be expressed by the following equation:

$$Re_p = \frac{\rho U D_{RBC}}{\mu} \quad (2)$$

where  $U$  is the average velocity of the RBCs flowing at the centerline of the hyperbolic microchannel,  $D_{RBC}$  is the undeformed diameter of a human RBC [1],  $\mu$  and  $\rho$  are the shear viscosity and density of Dextran 40 solution, respectively. It is interesting to note that the DI values initially increase monotonically with  $Re_p$ , but start converging for  $Re_p > 1$ , possibly indicating the existence of a maximum value above which the cells could not deform any further.

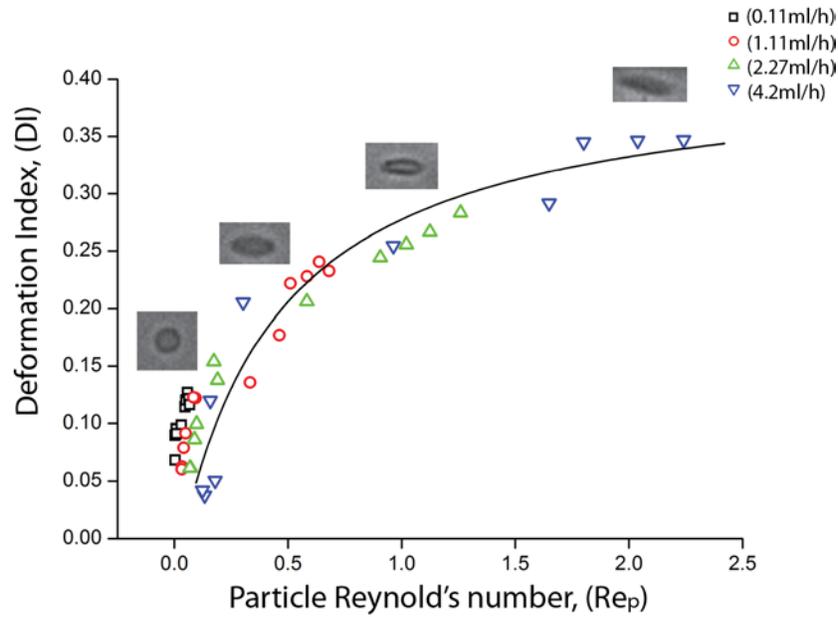


Figure 8. RBCs DI as a function of  $Re_p$  at the centerline of the hyperbolic microchannel for different volumetric flow rates:  $Q = 0.11\text{ml/h}$ ,  $Q = 1.11\text{ml/h}$ ,  $Q = 2.27\text{ml/h}$  and  $Q = 4.2\text{ml/h}$ . The  $Re_p$  was calculated by assuming the average diameter of a human RBC ( $D_{RBC} = 8\ \mu\text{m}$ ) [1].

### 3.3 Cell free layer

Over the past few years, many investigations have been conducted to improve the understanding of fluid dynamics and physics about the flow of rigid suspended particles. Although much is known for the flow of rigid suspended particles, much less is known when the suspended particles are deformable such as human RBCs. To the best of our knowledge, we report for the first time the effect of the flow rate on spatial distribution of the deformable human RBCs downstream of a hyperbolic-shaped contraction. Our results are compared with past results [10, 11] obtained in microchannels having sudden contractions.

The widths of cell depleted region, also known as cell free layer (CFL), as well as the widths of the core RBC region were measured upstream and downstream of the hyperbolic contraction region as shown in Figure 9.

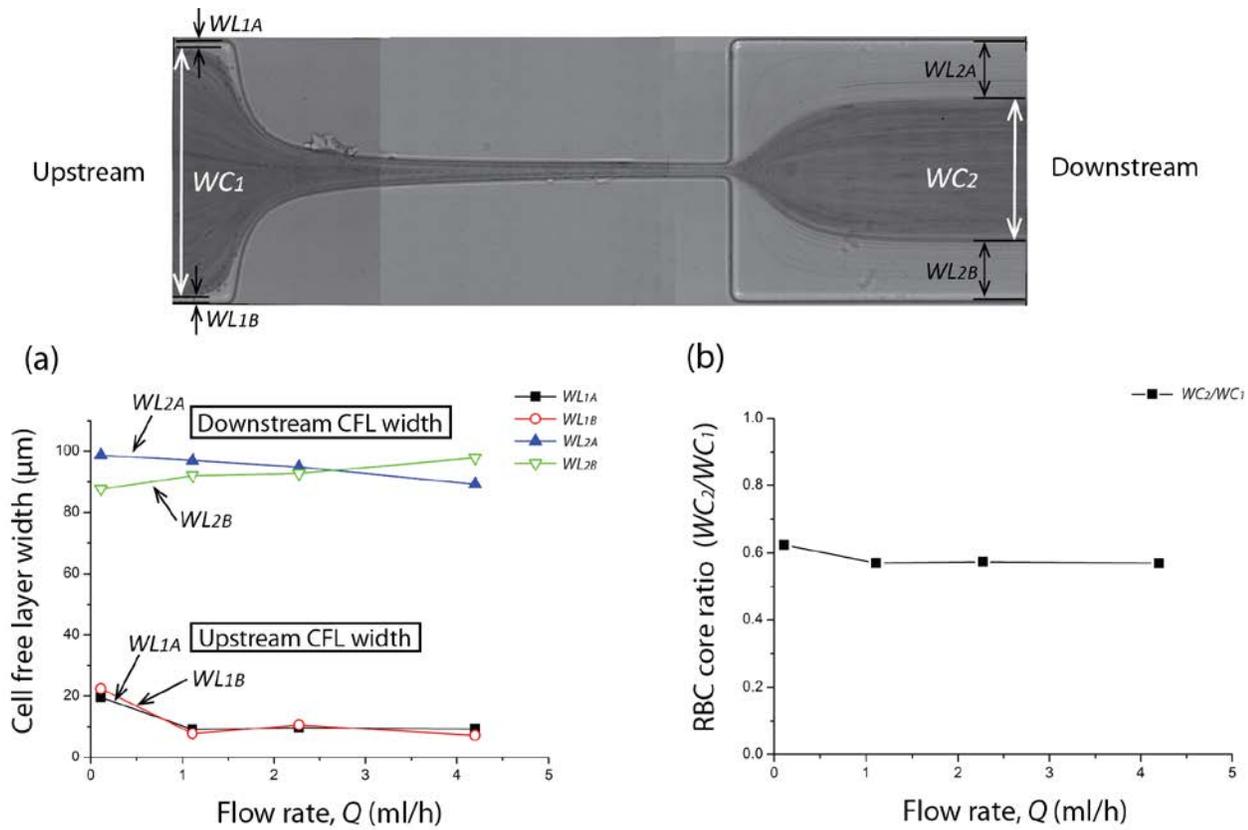


Figure 9 Cell free layer measurements. (a) Width of CFLs upstream and downstream as a function of the flow rate:  $Q = 0.11$  ml/h,  $Q = 1.11$  ml/h,  $Q = 2.27$  ml/h and  $Q = 4.2$  ml/h. (b) Effect of the flow rate on the ratio of the core RBC widths downstream and upstream of the contraction.

It is clear from the results shown in Figure 9(a) that the flow is highly symmetric both upstream and downstream of the contraction region, with the widths of the cell free layers being approximately the same on each side of the  $y = 0$  axis (corresponding to subscripts A and B). More importantly, it was observed for all the flow rates that the CFLs downstream of the expansion plane are greatly enhanced (by nearly ten times except for the lowest flow rate) relative to the CFLs upstream of contraction region. The enhancement of the CFLs downstream of contraction-expansion regions has been reported previously for abrupt contractions [3, 11] and as suggested this effect and geometry could be useful for continuous separation of RBCs from plasma. Compared with the experimental results performed in microchannels having sudden contractions [11], our results show higher enhancement of the downstream CFL, but a direct comparison is difficult given the different geometrical dimensions, such as the aspect ratio and expansion ratio. Further detailed investigations by studying more similar geometrical and physiological parameters need to be conducted in order to elucidate the true effect of *hyperbolic* contractions on the enhancement of the downstream CFL. Another important observation is that the widths of the CFLs and RBC core do not seem to vary significantly with flow rate. This is also clear in Figure 9(b), where we present the ratio of RBC core widths upstream ( $WC_1$ ) and downstream ( $WC_2$ ) of the contraction region as a function of the flow rate. A slightly higher ratio is observed only for the lowest flow rate due to the wider upstream CFL. A possible explanation for this phenomenon is that the CFL tends to increase both with the strain rate and with the time that RBCs spend within the hyperbolic contraction, and as a result by increasing the flow rate these two effects tend to cancel. Additionally, our results corroborate the study performed by Faivre et al. [11] where by using a sudden contraction-expansion they showed that not only

that the CFL is enhanced by the presence of a contraction but also that the flow rate has a weak effect on the thickness of the downstream CFL, for the range of flow rates and Reynolds tested.

#### 4 CONCLUSIONS, POTENTIAL BENEFITS AND FUTURE DIRECTIONS

The understanding of RBC deformability has become important ever since the clinical relevance of this hemorheological parameter has been recognized. The use of microfluidic devices, as proposed in this study, has the great advantage of being suitable to deal with single cell deformability while probing large numbers of cells in one single run. This high throughput capability together with the ability to achieve a controlled flow, make it possible to detect small changes of RBC deformability potentially in a more efficient and less time-consuming way than in conventional deformability measurements, having the potential to become a valuable approach for routine cell screening and early detection of several diseases related to RBC deformability. In this study, we focus on the effect of extensional flow, rather than shear flow, using a microfluidic device with a hyperbolic-shaped contraction in which the RBCs experience a strong extensional flow with a region of homogeneous strain rate along the centerline, as shown by our velocity measurements using particle tracking. Our measurements show that the RBCs are highly deformable under strong extensional flows, with the degree of deformability being highly dependent on the flow rate and exhibiting a DI that seems to be converging to a maximum plateau value as  $Re_p$  is increased. For the range of conditions tested here, RBC rupture was never observed. The proposed strategy to assess RBC deformability can be easily transformed into a simple, inexpensive and convenient tool suitable to detect and diagnose RBC related diseases. However, the viability to convert the proposed microfluidic device into a viable clinical instrument and the potential clinical significance of our measurements still requires further research.

Furthermore, there is a growing interest to develop microfluidic strategies that are able to perform deformation and plasma separation in order to be integrated into a single point-of-care (POC) diagnostic system, with no time-consuming pre-handling steps for the separation and labeling of cells. Our quantitative results clearly show that for the flow rate ranges covered in this work, the cell-free layer (CFL) thickness downstream of the expansion plane increases significantly (up to ten times), similarly to what is seen in abrupt contractions [3, 10, 11]. Consequently, this approach can be used as a microfluidic component for continuous separation of cells from plasma. Although our results indicate that the homogeneous extensional flow field generated in the hyperbolic contraction followed by the sudden expansion flow promote high CFL enhancement, further research with different geometrical and physiological parameters is needed, particularly with respect to the effect of the Hct, the channel aspect ratio, the expansion ratio, and the length of the contraction region. Another possible potential application, subject of future research, will be the sorting of blood cells, such as white cells, platelets and RBCs which differ in size and deformability. A recent work performed by Park and Jung [28] has shown that by using microchannels with a series of sudden contraction/expansions it is possible to separate rigid microspheres with different sizes. Hence, it is possible that by using a series of hyperbolic microchannels we may be able to develop a microfluidic device capable of sorting not only different kinds of blood cells but also cells with different degree of deformability.

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## References

- <sup>1</sup> C. Caro, T. Pedley, R. Schroter, and W. Seed, *The mechanics of the circulation* (Oxford University Press, 1978)
- <sup>2</sup> R. Skalak and P-I Branemark, *Science* **164**, 717-719 (1969).
- <sup>3</sup> M. Abkarian, M. Faivre, R. Horton, K. Smistrup, C. A. Best-Popescu, and H. A. Stone, *Biomed. Mater.* **3**, 034011 (2008).
- <sup>4</sup> M. R. Hardeman and C. Ince, *Clin. Hemorheol. Micro.* **21**, 277–284 (1999).
- <sup>5</sup> Y. I. Cho, M. P. Mooney, and D. J. Cho, *J. Diabetes. Sci. Technol.* **2**, 1130- 1138 (2008).
- <sup>6</sup> H. W. Hou, A. A. Bhagat, A. G. Chong, P. Mao, K. S. Tan, J. Han and C. T. Lim, **10**, 2605–2613 (2010)
- <sup>7</sup> K. Tsukada, E. Sekizuka, C Oshio, and H. Minamitani, *Microvasc. Res.* **61**, 231-239 (2001).
- <sup>8</sup> S. Shin, Y. Ku, J. Ho, Y. Kim, J. Suh, and M. Singh, *Clin. Hemorheol. Micro.* **36**, 253–261 (2007).
- <sup>9</sup> D.R. Gossett, H. T. K. Tse, S. A. Lee, Y. Ying, A. G. Lindgren, O. O. Yang, J. Rao, A. T. Clark, and D. Di Carlo. *Proceedings of the National Academy of Sciences of United States of America*, **109** (20), pp. 7630–7635 (2012).
- <sup>10</sup> M. Abkarian, M. Faivre, R. Horton, K. Smistrup, C. A. Best-Popescu, and H. A. Stone, *Biomed. Mater.* **3**, 034011 (2008).
- <sup>11</sup> M. Faivre, M. Abkarian, K. Bickraj and H. A. Stone. *Biorheology* **43** (2), 147 (2006).
- <sup>12</sup> V. Leble, R. Lima, R. P. Dias, C. S. Fernandes, T. Ishikawa, Y. Imai, and T. Yamaguchi, *Biomicrofluidics* **5**, 044120 (2011).
- <sup>13</sup> H. Fujiwara, T. Ishikawa, R. Lima, N. Matsuki, Y. Imai, H. Kaji, M. Nishizawa, and T. Yamaguchi, *J. Biomech.* **42**, 838 (2009).
- <sup>14</sup> S. Yang, A. Undar, and J. D. Zahn, *Lab Chip* **6**, 914 (2006).
- <sup>15</sup> S. S. Shevkoplyas, T. Yoshida, L. L. Munn, and M. W. Bitensky, *Anal. Chem.* **77**, 933 (2005).
- <sup>16</sup> S. S. Shevkoplyas, T. Yoshida, S. C. Gifford, and M. W. Bitensky, *Lab Chip* **6**, 914 (2006).
- <sup>17</sup> M. Musielak. *Clin. Hemorheol. Micro.* **42**, 47–64 (2009).
- <sup>18</sup> T. Shiga, N. Maeda, and K. Kon, *Crit. Rev. Oncol. Hemat.* **10** (1) 9-48 (1990).
- <sup>19</sup> D. H. Kim, P. K. Wong, J. Park, A. Levchenko and Y. Sun, *Annu. Rev. Biomed. Eng.* **11**, 203 (2009).
- <sup>20</sup> G. Tomaiuolo, M. Barra, V. Preziosi, A. Cassinese, B. Rotoli and S. Guido, *Lab Chip*, **11**, 449–454 (2011).
- <sup>21</sup> R. Zhao, J. F. Antaki, T. Naik, T. N. Bachman, M. V. Kameneva, and Z. J. Wu, *Biorheology* **43** (6) 747-65 (2006).
- <sup>22</sup> M. Kang, H. S. Ji, and K. C. Kim, *Int. J. Biol. Bio-med. Eng.* **1** (2) 1-8 (2008).
- <sup>23</sup> H. A. Barnes, *A handbook of elementary rheology*, (The University of Wales Institute of Non-Newtonian Fluid Mechanics, 2000).
- <sup>24</sup> D. F. James, G. M. Chandler, and S. J. Arnour, *J. Non-Newton. Fluid.* **35**, 421-443 (1990).
- <sup>25</sup> P. C. Sousa, F. T. Pinho, M. S. N. Oliveira, and M. Alves, *Biomicrofluidics* **5**, 014108 (2011).
- <sup>26</sup> S. S. Lee, Y. Yim, K. H. Ahn, and S. J. Lee, *Biomed. Microdevices.* **11**, 1021 (2009).
- <sup>27</sup> M. S. N. Oliveira, M. A. Alves, F. T. Pinho, and G. H. McKinley, *Exp. Fluids.* **43**, 437-451 (2007).
- <sup>28</sup> J. S. Park and H. I. Jung *Anal. Chem.* **81**(20), 8280-8288 (2009).
- <sup>29</sup> M. D. Abramoff, P. J. Magalhães, and S. J. Ram, *Biophotonics International* **11**, 36-42 (2004).
- <sup>30</sup> R. Lima, M. S. N. Oliveira, T. Ishikawa, H. Kaji, S. Tanaka, M. Nishizawa, and T. Yamaguchi, *Biofabrication* **1**, 035005 (2009).
- <sup>31</sup> R. Lima, T. Ishikawa, Y. Imai, M. Takeda, S. Wada, and T. Yamaguchi, *Ann. Biomed. Eng.* **37**, 1546 (2009).
- <sup>32</sup> E. Meijering, I. Smal, and G. Danuser, *IEEE Signal Proc. Mag.* **3**, 46 (2006).
- <sup>33</sup> T. Omori, T. Ishikawa, Y. Imai, and T. Yamaguchi, *J Biomech.* **46** (3) 548-553 (2013).
- <sup>34</sup> T. Omori, Y. Imai, T. Yamaguchi, and T. Ishikawa, *Phys. Rev. Lett.* **108** 138102 (2012).
- <sup>35</sup> E. A. Evans and Y.C. Fung, *Microvasc. Res.* **4** 335-347 (1972).
- <sup>36</sup> R. M. Hochmuth and R.E. Waugh, *Annu. Rev. Physiol.* **49** 209-219 (1987).
- <sup>37</sup> S. R. Keller and R. Skalak, *J. Fluid. Mech.* **120**, 27-47 (1982).
- <sup>38</sup> J. G. G. Dobbe, M. R. Hardeman, G. J: Streekstra, J. Strackee, C. Ince, and C. A. Grimbergen, *Blood. Cells. Mol. Dis.* **28**, 373-384 (2002).
- <sup>39</sup> G. G. Dobbe, M. R. Hardeman, G. J: Streekstra, and C. A. Grimbergen, *Biorheology*, **41**, 65-77 (2004).