Modelling wave–current interactions off the east coast of Scotland

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Abstract. Densely populated coastal areas of the North Sea are particularly vulnerable to severe wave conditions, which overtop or damage sea defences leading to dangerous flooding. Around the shallow southern North Sea, where the coastal margin is lying low and population density is high, oceanographic modelling has helped to develop forecasting systems to predict flood risk. However, coastal areas of the deeper northern North Sea are also subject to regular storm damage, but there has been little or no effort to develop coastal wave models for these waters. Here, we present a high spatial resolution model of northeast Scottish coastal waters, simulating waves and the effect of tidal currents on wave propagation, driven by global ocean tides, far-field wave conditions, and local air pressure and wind stress. We show that the wave–current interactions and wave–wave interactions are particularly important for simulating the wave conditions close to the coast at various locations. The model can simulate the extreme conditions experienced when high (spring) tides are combined with sea-level surges and large Atlantic swell. Such a combination of extremes represents a high risk for damaging conditions along the Scottish coast.

1 Introduction

Due to its semi-enclosed morphology and shoaling bathymetry, the North Sea experiences extreme wave conditions, in particular during winter periods (Woelf et al., 2002). When combined with sea-level surges such events can lead to damaging inundation of low-lying coastal regions, due to wave overtopping of sea defences. Development of a modelling and predictive capability for high-resolution wave conditions in the North Sea is therefore a high priority. However, the task is complicated due to interaction between locally generated waves and incoming swell from outside the region, and especially due to interactions between waves and tidal currents.

Crossing, or bimodal, sea states occur between 5 and 40 % of the time in the North Sea (Guedes Soares, 1984). These are generated when swell waves propagating into the region from distant storm events interact with locally generated waves which may be of very different direction, period, and height. Swell waves from the North Atlantic and the Norwegian Sea propagate into the North Sea, interacting with local wind-sea-generated waves, modifying the main spectral parameters. The interaction between differing wave trains is not fully understood, but crossing seas have been statistically associated with freak wave incidence, and shipping accidents (Waseda et al., 2011; Tamura et al., 2009; Cavaleri et al., 2012; Onorato et al., 2006, 2010; Sabatino and Serio, 2015; Toffoli et al., 2011a). The North Sea is particularly prone to rogue wave events, such as the famous Draupner wave recorded in 1994 (Haver, 2004), the first ever recorded rogue wave event, that occurred in crossing sea conditions (Adcock et al., 2011). In the paper, the effect of the enhancement of the wind-sea waves due to swell is assessed during storms.

In addition to crossing seas, wave–current interactions are a well-known cause of wave height amplification or attenuation. Wave–current interactions (WCIs) are depth- and current-induced modification of wave features. A seminal
study carried out by Tolman (1991) highlighted that WCIs are significant in the North Sea, changing the significant wave height \((H_s)\) and the mean wave period \((T_m)\) by 5 and 10\%, respectively, during storm periods. However, the model that was used by Tolman (1991) to assess this effect was at very coarse resolution and broad scale. In particular, the effect of the WCI in the coastal shallow areas was not considered. Phillips (1977) showed that in the absence of wave breaking, local wave amplitude is given by

\[
\frac{A}{A_0} = \frac{c_0}{\sqrt{c(c + 2U)}},
\]

where \(A\) is the resulting wave amplitude, \(A_0\) is the unperturbed amplitude of the wave field, \(c\) is the wave phase speed, and \(U\) is the current that interacts with the wave train. It is important to notice that the sign of the current is determinant on the effect of the WCI: if the current travels in opposite direction with the wave train, there will be an enhancement of the significant wave height. Conversely, if current and waves are in the same direction, the \(H_s\) decreases. For deep water waves, the phase speed \(c\) depends only on the period of the wave, while in shallow water \(c\) depends only on the depth. Equation (1) shows that the waves travelling in a direction opposing the current \((U < 0)\) have a positive ratio \(A/A_0\) and consequently, an enhancement of the wave amplitude.

WCIs could also lead to the breaking of the wave: if the current is strong enough to block the wave train (Ris and Holthuijsen, 1996), these waves can break and lose energy before arriving to the coastline (Chawla and Kirby, 2002, 1998).

WCIs are particularly difficult to quantify empirically, and computationally intensive to model. In addition, the WCIs are a well-known mechanism for the formation of rogue waves in the ocean: the Agulhas current, that flows near the coastline of South Africa, was one of the first places in which this mechanism was identified (Mallory, 1974; Lavrenov, 1998; Lavrenov and Porubov, 2006). Recently, many studies (Onorato et al., 2011; Toffoli et al., 2013, 2015; Shirira and Slunyaev, 2014; Ma et al., 2013) highlighted that the interaction between a train of waves with an adverse current could increase its wave steepness, and cause rogue waves due to the modulational instability (Benjamin and Feir, 1967). However, it is clear that shallow coastal waters, embayments, and headlands are particular foci for interactions (Hearn et al., 1987; Signell et al., 1990a). Model studies have concentrated on comparing wave height and period in coupled and uncoupled model versions showing, for example, 3 % difference in wave height and 20 % in wave period in the Dutch and German coastal waters of the North Sea (Osuna and Monbaliu, 2004). Similar results have been obtained for coastal waters of the Adriatic during bora conditions (Benetazzo et al., 2013), finding a maximum reduction for the \(H_s\) of 0.6 m in the central Adriatic and a simultaneous increase up to 0.5 m in the gulfs of Trieste and Venice. WCIs were also studied during hurricane conditions off the eastern seaboard of the USA (Xie et al., 2008).

Sea defences of coastal settlements along the northeast coast of Scotland have suffered several damaging events during the period 2009–2014 as a result of surges and waves. The coastal waters are dominated by strong tidal currents and wind-driven residuals, are exposed to wave trains entering the North Sea from the north, and generated by storm events in the central and southern North Sea. Although the oceanography of the North Sea as a whole has been intensively studied since the 1830s (Whewell, 1830; Proudman and Doodson, 1924; Dietrich, 1950; Huthnance, 1991; Otto et al., 1990), and the region was one of the earliest to be subjected to computational hydrodynamic modelling (Flather, 1987; Davies et al., 1985), high-resolution modelling activity has been largely concentrated in areas with potential for wave and tidal energy extraction (Adcock et al., 2013; Bryden and Couch, 2006; Baston and Harris, 2011; Shields et al., 2011, 2009). However, there are no such models for the northeast coast of mainland Scotland, and none which include coupled WCIs. Our objective here was to develop and test such a model for the stretch of coastline between the Firth of Tay and Peterhead, centred on the strategically important port city of Aberdeen and the town of Stonehaven (Fig. 1). The latter is the base for a governmentally supported marine monitoring site with a >15-year time series of high-resolution data on a wide range of environmental parameters (Bresnan et al., 2009).

2 Materials and methods

The MIKE by DHI model was used to simulate the tidal- and wind-driven circulation, and the wave propagation. The MIKE software is composed of different modules, for the creation of a model grid and input files and to simulate different hydrodynamical features at the same time or separately. The following modules were used:

- The MIKE Zero modules for generating the computational grid and the input files.
- The MIKE 3 FM module for simulating the tidal and the wind-driven circulation.
- The MIKE 21 SW module for modelling the wave propagation.

For the simulation of the WCIs, a one-way coupling between MIKE 3 FM and MIKE 21 SW was set up. The depth average flow fields and the water level output from MIKE 3 FM was provided as input to the MIKE 21 SW model.

2.1 The computational grid

Both MIKE 3 FM and MIKE 21 SW use an unstructured grid approach, with triangular elements (Ferziger and Perić,
Unstructured grids can represent complex coastlines better than a rectangular grid and potentially provide more realistic flows, enabling the geography of the coastline to affect the propagation of tidal and surface waves in a realistic manner. In addition, triangular grid elements allow smoothly changing cell sizes across a region, with the highest resolution concentrated in an area of particular interest. The mesh for the area of study is shown in Fig. 2. An enhanced-resolution area was created near Stonehaven, because this work is part of a wider study focusing on the resuspension of the sediments in this part of the domain. This high-resolution area also covered the Firth of Forth and the Aberdeenshire coastline, where previous studies have shown enhanced currents due to interaction between the tidal wave and the Scottish coastline (Dietrich, 1950; Otto et al., 1990).

2.2 The MIKE 3 FM hydrodynamic model

MIKE 3 FM (flow model) is based on the numerical solution of the 3-D incompressible Reynolds-averaged Navier–Stokes equations, under the Boussinesq and the hydrostatic pressure approximations (DHI, 2011a). The spatial discretization of the primitive equations is performed using a cell-centred finite-volume method. In the finite-volume method the volume integrals in the partial differential equations with a divergence are converted to surface integrals using the Gauss–Ostrogradsky theorem (Toro, 2009).

MIKE 3 FM has a flexible approach for simulating the flow in the water column. It is possible to choose between sigma layers (Song and Haidvogel, 1994), $z$ layers, and coupled sigma and $z$ layers. For our purpose we decided to use the equidistant sigma layers approach, because the bathymetry of the area was not sufficiently complex to require a more accurate description with a coupled sigma and $z$ layer that would be extremely computationally expensive. Sigma layers are also useful for resolving the water column well throughout the tidal cycle, given the large tidal range.
The coupled sigma and $z$ layers were tested, but there was no significant improvement for simulating the flow.

For the horizontal eddy viscosity the formulation proposed by Smagorinsky (1963) was used, in which the sub-grid-scale transport is expressed by an effective eddy viscosity related to a characteristic length scale rather than a constant eddy viscosity. This sub-grid-scale viscosity is given by

$$A = c_s^2 l^2 \sqrt{2 S_{ij} S_{ij}},$$  \hspace{1cm} (2)$$

where $c_s$ is the Smagorinsky constant, $l$ is the characteristic length of the grid size, and the deformation rate is given by (Lilly, 1966; Deardorff, 1971; Smagorinsky, 1963)

$$S_{ij} = \frac{1}{2} \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right).$$  \hspace{1cm} (3)$$

The bed resistance was parameterized using a constant quadratic drag coefficient $c_f$. The average bottom stress is determined by a quadratic friction law:

$$\tilde{\tau}_b = c_f \rho_0 u_b |\tilde{u}_b|,$$  \hspace{1cm} (4)$$

where $\tilde{u}_b$ is the average flow velocity above the bottom and $\rho_0$ is the density of the water. The value of the above parameters chosen for the model are reported in Sect. 3.1.

The model was forced with a time series of tidal elevations at the open boundaries from the open-source OSU (Oregon State University) Tidal Prediction Software (OTPS) (Egbert et al., 2010), based on TOPEX satellite observation of the water-level observations interpolated with tide gauge data from the European shelf region. In order to take account of the wind-driven circulation and surge in the model, meteorological forcing was applied across the model domain, using the ERA-Interim reanalysis for wind velocity and mean sea-level pressure (Dee et al., 2011).

### 2.3 The MIKE 21 SW wave model

The MIKE 21 SW (spectral wave) is an unstructured grid model for wave prediction and analysis (DHI, 2011b). The MIKE 21 SW is based on the wave action conservation equation (Komen et al., 1996; Young, 1999), where the dependent variable is the frequency-directional wave action spectrum. This is given by

$$\frac{\partial N}{\partial t} + \nabla \cdot (c N) = \frac{S}{\sigma},$$  \hspace{1cm} (5)$$

where $S$ is the energy source term, defined as

$$S = S_m + S_{nl} + S_{ds} + S_{bot} + S_{surf}$$  \hspace{1cm} (6)$$

that depends on the energy transfer from the wind to the wave field $S_m$, on the nonlinear wave–wave interaction $S_{nl}$, on the dissipation due to depth-induced wave breaking $S_{surf}$, on the dissipation due to bottom friction $S_{bot}$, and on the dissipation caused by the white-capping $S_{ds}$.

The wave action density spectrum $N(\sigma, \theta)$ is defined as (Bretherton and Garrett, 1968)

$$N(\sigma, \theta) = \frac{E(\sigma, \theta)}{\sigma},$$  \hspace{1cm} (7)$$

where $E$ is the wave energy density spectrum, $\sigma = 2\pi f$ is the angular frequency (where $f$ is the frequency), and $\theta$ is the direction of wave propagation. The momentum transfer from the wind to the waves follows the formulation in Komen et al. (1996). The momentum transfer and the drag depend not only on the strength of the wind but also on the wave state itself.

For the physics of the propagation and breaking of the waves we choose the following parameters:

- The depth-induced wave breaking is based on the formulation of Battjes and Janssen (1978), in which the gamma parameter is a constant 0.6 across the domain. The formulation of the depth-induced wave breaking can be written as

$$S_{surf}(\sigma, \theta) = -\frac{\alpha Q_b \bar{\sigma} H_m^2 E(\sigma, \theta)}{8\pi E_{tot}},$$  \hspace{1cm} (8)$$

where $\alpha \approx 1.0$ is a calibration constant, $Q_b$ is the fraction of breaking waves, $\bar{\sigma}$ is the spectrum average frequency, $E_{tot}$ is the total wave energy that is linked to the wave action density spectrum, and $H_m$ is the estimated maximum wave height, that is defined as $H_m = \gamma d$ (Battjes and Janssen, 1978), in which $d$ is the depth and $\gamma$ is the free breaking parameter (Battjes, 1974).

- The bottom friction is specified in the model as the Nikuradze roughness ($k_0$) (Nikuradse, 1933; Johnson and Kofoed-Hansen, 2000).

- The white-capping formulation described in Komen et al. (1996) in order to consider the dissipation of waves, based on the theory of Hasselmann (1974). For the fully spectral formulation, the white capping assumes a form that is dependent on the mean frequency $\bar{\sigma}$ and on the wavenumber $k$:

$$S_{ds}(\sigma, \theta) = -C_{ds}\left(\frac{k^2 m_0}{\bar{\sigma} N(\sigma, \theta)}\right)^2$$

$$\left[ (1 - \delta) \frac{k}{\bar{\sigma}} + \delta \left(\frac{k}{\bar{\sigma}}\right)^2 \right] \frac{\partial N}{\partial \sigma},$$  \hspace{1cm} (9)$$

Here the two parameters, $C_{ds}$ and $\delta$, are the two dissipation coefficients that control the overall dissipation rate and the strength of dissipation in the energy/action spectrum, respectively, and $m_0$ is the zeroth moment of the overall spectrum.

The values of the above parameters chosen for the model are reported in Sect. 3.2.
The model also included nonlinear energy transfer such as the quadruplet wave interaction (Komen et al., 1996) and the triad wave interaction which is the dominant nonlinear interaction in shallow water (Eldeberky and Battjes, 1995, 1996).

The forcings included in the model are the local wind and the swell wave field from outside the model area and specified at the model boundaries. For the model boundaries we used boundary conditions from the Venugopal and Nemalidinne (2014, 2015) North Atlantic model, a larger aries we used boundary conditions from the Venugopal and (1996). The ERA-Interim 0.125° × 0.125° model was used to provide a wind field across the model domain with a time resolution of 6 h (Dee et al., 2011; Berrisford et al., 2011).

2.4 Wave–current interactions

The WCIs are implemented using a one-way coupling between currents and waves. The model was run without and with currents implemented, and then the differences between the two runs were studied. The WCIs in the MIKE model are taken in account in the dispersion relation for the angular frequency term, since the current due by tides and wind affect the propagation and changes the wavelength of the wave train. The MIKE 21 SW dispersion relation in fact is

\[ \sigma = \sqrt{gk \tanh(kd)} = \omega - k \cdot U. \]

The one-way coupling has, however, some limitations since it is not taking into account the modification of the current by the wave itself (Michaud et al., 2011; Bennis et al., 2011).

2.5 Swell detection

In the present study, the wind-sea and the swell waves and their interaction are studied. MIKE 21 SW gives the opportunity to separate spectrally the windsea waves and the swell waves. There are two criteria, based on a dynamic threshold, available to make this separation.

The first criterion is based on the difference of the energy between the spectrum and the fully developed sea condition (Earle, 1984). In this case, the threshold frequency is identified as

\[ f_{\text{threshold}} = \alpha f_{p,\text{PM}} \left( \frac{E_{\text{PM}}}{E_{ \text{Model} }} \right)^\beta, \]

where \( \alpha = 0.7 \), \( \beta = 0.31 \), \( E_{\text{Model}} \) is the total energy at each node point calculated by the MIKE 21 SW model, and the Pierson–Moskowitz peak frequency and the energy are estimated as

\[ f_{p,\text{PM}} = 0.14 \frac{g}{U_{10}} \]

\[ E_{\text{PM}} = \left( \frac{U_{10}}{1.4g} \right)^4. \]

The validation for waves was carried out using four statistical indices: the bias, the root mean square error (RMSE), the correlation coefficient (R), and the scatter index (SI). These indices are defined below.

\[ \text{Bias} = \frac{1}{N} \sum_{i=1}^{N} (x_{oi} - x_{mi}) \]

\[ \text{RMSE} = \sqrt{\frac{1}{N} \sum_{i=1}^{N} (x_{oi} - x_{mi})^2} \]
Table 1. Location of the validation/calibration instrumentation.

<table>
<thead>
<tr>
<th>Description</th>
<th>Coordinates</th>
<th>Depth (m)</th>
<th>Use</th>
</tr>
</thead>
<tbody>
<tr>
<td>Aberdeen tide gauge</td>
<td>-2.0803</td>
<td>57.144</td>
<td>– water level val/cal</td>
</tr>
<tr>
<td>Leith tide gauge</td>
<td>-3.1682</td>
<td>55.9898</td>
<td>– water level validation</td>
</tr>
<tr>
<td>Buckie tide gauge</td>
<td>-2.9667</td>
<td>57.6667</td>
<td>– water level validation</td>
</tr>
<tr>
<td>Firth of Forth buoy</td>
<td>-2.5038</td>
<td>56.1882</td>
<td>– waves validation</td>
</tr>
<tr>
<td>Moray Firth buoy</td>
<td>-3.3331</td>
<td>57.9663</td>
<td>– waves val/cal</td>
</tr>
<tr>
<td>Aberdeen wave rider</td>
<td>-2.0500</td>
<td>57.1608</td>
<td>– waves validation</td>
</tr>
<tr>
<td>BODC 4551 RCM</td>
<td>-2.8000</td>
<td>57.7910</td>
<td>12 current validation</td>
</tr>
<tr>
<td>BODC 4561 RCM</td>
<td>-1.9680</td>
<td>57.2320</td>
<td>12 current validation</td>
</tr>
<tr>
<td>BODC 4562 RCM</td>
<td>-1.9680</td>
<td>57.2320</td>
<td>27 current validation</td>
</tr>
<tr>
<td>BODC 4571 RCM</td>
<td>-1.9020</td>
<td>57.2260</td>
<td>12 current validation</td>
</tr>
<tr>
<td>BODC 4572 RCM</td>
<td>-1.9020</td>
<td>57.2260</td>
<td>52 current validation</td>
</tr>
<tr>
<td>BODC 4582 RCM</td>
<td>-2.1500</td>
<td>56.9870</td>
<td>23 current validation</td>
</tr>
<tr>
<td>BODC 4591 RCM</td>
<td>-2.0980</td>
<td>56.9820</td>
<td>12 current validation</td>
</tr>
<tr>
<td>BODC 4592 RCM</td>
<td>-2.0980</td>
<td>56.9820</td>
<td>47 current validation</td>
</tr>
</tbody>
</table>

\[ R = \frac{1}{N} \sum_{i=1}^{N} (x_{o_i} - x_o) (x_{m_i} - x_m) \]  
\[ SI = \frac{\text{RMSE}}{\text{max}(x) - \text{min}(x)} \]

For tidal current validation we used, instead of the SI, the normalized root mean square error (NRMSE), that is defined as

\[ \text{NRMSE} = \frac{\text{RMSE}}{\text{max}(x) - \text{min}(x)} \]

Table 2. Computed RMSE for the main harmonic components, the validation for each tide gauge is reported in the Table S1.

<table>
<thead>
<tr>
<th>Components</th>
<th>RMSE</th>
</tr>
</thead>
<tbody>
<tr>
<td>$M_2$</td>
<td>2.52</td>
</tr>
<tr>
<td>$S_2$</td>
<td>1.64</td>
</tr>
<tr>
<td>$N_2$</td>
<td>1.31</td>
</tr>
<tr>
<td>$O_1$</td>
<td>0.76</td>
</tr>
<tr>
<td>$K_1$</td>
<td>0.44</td>
</tr>
<tr>
<td>$Q_1$</td>
<td>0.42</td>
</tr>
</tbody>
</table>

Table 3 shows that the model adequately represents the current speeds in the domain. The validation shows that the model slightly underestimates the current, however it can be noticed that the bias of the model was very low. The RMSE, except for one observation, does not exceed 15 % of the maximum speed.

For tidal current validation we used, instead of the SI, the normalized root mean square error (NRMSE), that is defined as

\[ \text{NRMSE} = \frac{\text{RMSE}}{\text{max}(x) - \text{min}(x)} \]

The validation was performed for different years. For the hydrodynamic model the agreement between modelled and observed water level was evaluated for the entire year 2007, while the currents were validated for 1992, where the rotor current meter observations were available. The wave model was validated for 2010 and 2008, where observations and boundary inputs were available.

3 Results

3.1 Calibration and validation of the hydrodynamic model

The hydrodynamic model was calibrated for the year 2007, based on the agreement with the recorded water level at the tide gauge in Aberdeen. The calibration parameters were the time step that was fixed at 1 s after an analysis of the Courant–Friedrichs–Lewy (CFL) conditions (higher time steps were investigated but the model was unstable); the Smagorinsky constant that was set to 0.2 (for values > 0.3 the model showed some blow-up); and the bottom roughness that was parameterized with the drag coefficient $c_f = 0.0025$. After calibration, the MIKE 3 tidal model was validated against harmonic components extracted from both observed and modelled data for water level. The agreement between modelled and observed currents was also investigated. RMSE for the amplitude of harmonic components was less than 1 % for all the cases, while the phase of the main semidiurnal component was well modelled. In particular for the dominant $M_2$ component, the phase error was very low and the amplitude was well modelled (see Tables 2 and S1 in the Supplement for more details). The validation results show that the modelled results are in good agreement with the recorded tidal amplitude and phase. The model was run for 1992 and measurements obtained from BODC from eight locations were used to validate the currents in the model. The validation of the single components $u$ and $v$ is reported in Table S2. Table 3 shows that the model adequately represents the current speeds in the domain. The validation shows that the model slightly underestimates the current, however it can be noticed that the bias of the model was very low. The RMSE, except for one observation, does not exceed 15 % of the maximum speed.
Table 3. Results from the validation of the currents, showing the difference between the modelled and observed current speeds at the eight locations reported in Table 1.

<table>
<thead>
<tr>
<th>RCM no.</th>
<th>Lat</th>
<th>Long</th>
<th>Depth (m)</th>
<th>RMSE  (m s(^{-1}))</th>
<th>NRMSE (m s(^{-1}))</th>
<th>(R^2)</th>
<th>Bias (m s(^{-1}))</th>
</tr>
</thead>
<tbody>
<tr>
<td>4551</td>
<td>−2.8</td>
<td>57.791</td>
<td>12</td>
<td>0.094</td>
<td>0.157</td>
<td>0.17</td>
<td>0.001</td>
</tr>
<tr>
<td>4561</td>
<td>−1.968</td>
<td>57.232</td>
<td>12</td>
<td>0.111</td>
<td>0.124</td>
<td>0.70</td>
<td>−0.03</td>
</tr>
<tr>
<td>4562</td>
<td>−1.968</td>
<td>57.232</td>
<td>27</td>
<td>0.075</td>
<td>0.105</td>
<td>0.75</td>
<td>−0.02</td>
</tr>
<tr>
<td>4571</td>
<td>−1.902</td>
<td>57.226</td>
<td>12</td>
<td>0.223</td>
<td>0.147</td>
<td>0.23</td>
<td>−0.05</td>
</tr>
<tr>
<td>4572</td>
<td>−1.902</td>
<td>57.226</td>
<td>52</td>
<td>0.087</td>
<td>0.112</td>
<td>0.80</td>
<td>−0.02</td>
</tr>
<tr>
<td>4582</td>
<td>−2.15</td>
<td>56.987</td>
<td>23</td>
<td>0.075</td>
<td>0.124</td>
<td>0.80</td>
<td>0.02</td>
</tr>
<tr>
<td>4591</td>
<td>−2.098</td>
<td>56.982</td>
<td>12</td>
<td>0.125</td>
<td>0.132</td>
<td>0.73</td>
<td>−0.062</td>
</tr>
<tr>
<td>4592</td>
<td>−2.098</td>
<td>56.982</td>
<td>47</td>
<td>0.073</td>
<td>0.121</td>
<td>0.82</td>
<td>−0.05</td>
</tr>
</tbody>
</table>

3.2 Calibration and validation of the wave model

The calibration of the wave model was carried out for 3 months in 2008 and was based on the agreement between the observed and the modelled \(H_s\) of the Firth of Forth wave gauge. There were three calibration parameters: the wave-breaking parameter \(\gamma\) and the two dissipation coefficients associated to the wave breaking (\(C_{\text{dis}}\) and \(\delta\)). During this procedure we noticed how the most sensible parameter was the \(\gamma\) controlling the wave breaking. We investigated the behaviour of the \(\gamma\) for the range 0.6–1.0, since most of the observation studies in literature were reporting such values (Battjes, 1974; Stive, 1985; Battjes and Stive, 1985; Nelson, 1987, 1994; Kaminsky and Kraus, 1993), and we found that \(\gamma = 0.6\) was the value giving better results for the \(H_s\). The wave-breaking dissipation coefficients were fixed to \(C_{\text{dis}} = 2.5\) and \(\delta = 0.8\), while the Nikuradze bottom roughness was fixed to 0.01 m. The results of the wave gauges and satellite validation are reported in Table 4. We evaluated the performance of the wave model with WCIs implemented (coupled) and without WCIs (uncoupled). There was a good agreement between the modelled-with-WCI and measured wave data. The bias does not exceed 0.15 m for significant wave height. Table 4 and Fig. 3 shows that the model estimates correctly the significant wave height in the Firth of Forth and in Aberdeen but underestimates this parameter in the Moray Firth. However, the agreement with the data is still satisfactory. In particular, low RMSE values were recorded for the Aberdeen wave gauge, which is the only coastal shallow-water wave gauge that is available in the area (the depth of the mooring site is 10 m). The model performance against satellite data randomly sampled throughout the domain shows good agreement. Without the WCI included in the model, small or no differences were estimated for significant wave height, but larger differences were seen for mean wave period (\(T_{m01} = 2\pi m_0/m_1\)): the calculated RMSE for the uncoupled model was 0.97 s in Aberdeen, 1.24 s in the Firth of Forth and 1.83 s in the Moray Firth. Comparing satellite observations in spring and winter conditions, it is possible to conclude that, in general, the model provides accurate predictions for wave heights < 1.5–2 m, but slightly underestimates the height of larger waves. On the other hand, wave periods are better modelled in the winter period when the waves are higher. No or very small differences were recorded between the coupled and uncoupled models for satellite validation. This is because the resolution of the satellite data is low and because the satellite data are often in deep water, where the WCI are less important.

3.3 Wave–current interaction

Predicted wave field with and without WCI were compared during a 7-month period in 2010, covering both winter and summer conditions, for evaluating the importance of WCI on wave features. The results are shown in Fig. 4. For the comparison between the coupled and the uncoupled model the root mean square (rms) between the two runs was computed.
Table 4. Comparison between observed and modelled (both with and without WCIs implemented) wave heights and periods for wave gauges and satellite observations. The reported validation was carried out for 2010 (Firth of Forth and Moray Firth) and for 2008 (Aberdeen wave gauge and satellite observations). Details of the observation data are reported in Table 1 of the paper and in Table S2 of the Supplement.

<table>
<thead>
<tr>
<th>Location</th>
<th>Coupled Bias</th>
<th>Coupled RMSE</th>
<th>Coupled R</th>
<th>Coupled SI</th>
<th>Uncoupled Bias</th>
<th>Uncoupled RMSE</th>
<th>Uncoupled R</th>
<th>Uncoupled SI</th>
</tr>
</thead>
<tbody>
<tr>
<td>Firth of Forth</td>
<td>-0.02 m</td>
<td>0.30 m</td>
<td>0.941</td>
<td>0.27</td>
<td>-0.01 m</td>
<td>0.30 m</td>
<td>0.939</td>
<td>0.27</td>
</tr>
<tr>
<td></td>
<td>-0.70 s</td>
<td>1.17 s</td>
<td>0.767</td>
<td>0.25</td>
<td>-0.76 s</td>
<td>1.24 s</td>
<td>0.758</td>
<td>0.27</td>
</tr>
<tr>
<td>Moray Firth</td>
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<td>0.42 m</td>
<td>0.849</td>
<td>0.38</td>
<td>-0.15 m</td>
<td>0.42 m</td>
<td>0.848</td>
<td>0.39</td>
</tr>
<tr>
<td></td>
<td>-1.18 s</td>
<td>1.75 s</td>
<td>0.668</td>
<td>0.39</td>
<td>-1.23 s</td>
<td>1.83 s</td>
<td>0.656</td>
<td>0.41</td>
</tr>
<tr>
<td>Aberdeen</td>
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<td>0.21 m</td>
<td>0.836</td>
<td>0.32</td>
<td>-0.07 m</td>
<td>0.22 m</td>
<td>0.831</td>
<td>0.32</td>
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<tr>
<td></td>
<td>-0.25 s</td>
<td>0.91 s</td>
<td>0.715</td>
<td>0.20</td>
<td>-0.30 s</td>
<td>0.97 s</td>
<td>0.701</td>
<td>0.21</td>
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<tr>
<td>Satellite</td>
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<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Winter</td>
<td>-0.2 m</td>
<td>0.4 m</td>
<td>–</td>
<td>0.25</td>
<td>-0.2 m</td>
<td>0.4 m</td>
<td>–</td>
<td>0.25</td>
</tr>
<tr>
<td></td>
<td>+0 s</td>
<td>0.8 s</td>
<td>–</td>
<td>0.15</td>
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<td>0.8 s</td>
<td>–</td>
<td>0.15</td>
</tr>
<tr>
<td>Spring</td>
<td>-0.1 m</td>
<td>0.3 m</td>
<td>–</td>
<td>0.21</td>
<td>-0.1 m</td>
<td>0.3 m</td>
<td>–</td>
<td>0.21</td>
</tr>
<tr>
<td></td>
<td>+0.1 s</td>
<td>1.2 s</td>
<td>–</td>
<td>0.23</td>
<td>+0.1 s</td>
<td>1.2 s</td>
<td>–</td>
<td>0.23</td>
</tr>
</tbody>
</table>

Figure 4. Root mean square difference between wave model output with and without WCI: (a) significant wave height (m), (b) peak wave period (s), (c) wave directional spreading (degrees).

Results show some differences between the two runs. In particular, the largest deviations due to WCI are found in coastal areas, such as around headlands and bays, and in estuaries, in which the currents (mostly driven by tides) are strongest. As expected, the highest differences were seen in the proximity of the coastline (Signell et al., 1990b): this was because the strength of the mainly tidal-driven currents are stronger (Dietrich, 1950; Otto et al., 1990). During spring tides, higher values for the current were recorded off northeast England and near Peterhead and Aberdeen (see Fig. 1). Wave periods are more affected than wave heights in this coupling, with rms deviations that can be on average 20% (absolute value) in shallow-water coastal areas. We also considered the effect of the WCIs on the wave directional spreading, as this is an important variable for the stability of the wave train in deep water and for its evolution (Benjamin and Feir, 1967). The results showed that during the 7-month period the significant wave height was, on average, less affected than directional spreading or wave periods: the difference was of the order of magnitude of 0.1 m near the coastline and less offshore, while the difference in peak spectral wave period ($T_p$) exceeded 1 s in some of the east coast firths such as the Moray Firth and the Firth of Forth.

Maximum positive and negative variation during the 7-month period were also studied (Figs. 5 and 6). The figure is similar to the RMSE: the larger variation is reported only in the coastal areas, while in the open sea the maximum variation is limited up to 1 m. Spatially, the maximum variation of $H_s$ between the coupled and the uncoupled run was $+2.8$ and $-1.8$ m, both occurring during storm events and both occurring in coastal areas, near the coast of Aberdeen and Pe-
terhead and south of the Firth of Forth, in which the tidally driven current is stronger (Dietrich, 1950; Otto et al., 1990).

3.4 Current and swell effect on the windsea wave field

In order to study the importance of the WCIs and the coupling between swell and windsea waves off the east coast of Scotland, three storms were considered in the period January–August 2010. Storm events were identified by examining the time series in the Firth of Forth and the Moray Firth in which the highest $H_s$ were recorded. These three storms were selected because they were the three most intense storms during the considered period and originated from different weather conditions.

3.4.1 The 26–27 February 2010 storm

Between 25 and 27 February 2010, the UK was affected by a low pressure system, that moved rapidly from west to east. From the afternoon of the 25th to the 26th, the centre of the storm was over the North Sea (Fig. 7). At the same time, another low pressure system (not shown in the map) was over the Norwegian Sea, causing a train of swell moving from north to south. Comparison of modelled and observation wave heights and wave period conditions for this storm are reported in Fig. 8. In addition, the modelled conditions in the Aberdeen wave rider location are reported. The figure shows that the model reproduces adequately the conditions during that storm, in particular around the time in which the maximum $H_s$ was reached.

The low pressure over North Sea caused windsea waves exceeding 4 m. In Fig. 9 the situation in the sea is shown at 12:30 UTC of the 26th: swell waves contributed to enhancing the $H_s$ in the centre of the storm, while a train of swell waves was forming from this storm, travelling west to the Moray Firth. Interaction of the windsea and the swell waves caused high waves along the east coast: the maximum recorded $H_s$ by the Firth of Forth wave gauge was 4.8 m. WCI contributed to the enhancement of $H_s$ by up to 1 m in coastal areas, while in the open sea the contribution was very low, up to 0.1 m. In the afternoon of the 26th (Fig. 10, at 19:00 UTC) the storm was near the Firth of Forth. The contribution of the swell waves was significant, increasing the $H_s$ by up to 1 m: model outputs showed that the central part of the storm had an $H_s > 5$ m, while without the swell coming from north the centre of the storm would have been an $H_s < 4.5$ m. To our knowledge, no significant damages were recorded for this storm.

3.4.2 The 30–31 March 2010 storm

The larger storm in 2010 occurred during the night of 30 March 2010. Between 29 March and 1 April 2010 the southeast coast of Scotland and the north of England were struck by severe weather and very strong winds. These conditions were caused by a strong depression that originated from a weak minimum near the Azores islands, in the North Atlantic, in front of the Portuguese coast. This low pressure was < 990 hPa once over Great Britain and Ireland at midnight of 30 March 2010 and reached its minimum the day after with a depression of < 980 hPa over the north of England. The evolution of the storm from surface pressure charts from ECMWF ERA-Interim reanalysis is reported in Fig. 11 (Dee et al., 2011; Berrisford et al., 2011). These figures clearly show that the depression, at its maximum strength, is just
above the south of Scotland during the night between 30 and 31 March 2010. This depression generated both very high waves ($H_s$ exceeded 6 m, measured in the Firth of Forth) and surge waves exceeding 0.5 m (measured both by the Aberdeen and Leith tide gauges). The waves caused significant damages to the coastal defences of cities in the southeast of Scotland. In particular, the City of Edinburgh Council estimated the damages to coastal defences to be about GBP 23 000. Also, in Berwick, at the southern entrance of the Firth of Forth, some damages were caused to the harbour infrastructures. To the east, in Dumbar, waves topped the roofs of two-floor houses.

Damaging conditions associated with this storm were caused by a combination of simultaneous factors: (1) tides in the spring period, (2) a surge wave of about 0.5 m generated by local pressure and wind, (3) windsea waves generated locally that interacted with strong currents, (4) a weak but significant swell waves field that interacted with the windsea waves.

Figure 12 shows the intensity of the current in the Aberdeen wave gauge location and the resulting WCI. It can be seen that the current was strongly enhanced by the wind, and consequently the WCI effect was stronger.

At about 00:30 UTC on 31 March 2010, the storm was at its maximum, causing the wave field to hit the coastline at around the same time as high tide and surge. The different components of the storm were analysed. First, the surge wave generated by the minimum of pressure above the North Sea was studied. Figure 13 shows the difference between the total water level and the water level due to tides at 02:00 UTC on 31 March 2015. The model predicted a surge wave up to 0.5 m. A comparison between the recorded water level and the model output showed that the model underestimated the surge wave by about 0.1 m. The reason for this underestimation could be because the boundary conditions for the model only included tidal water level and did not include the surge wave from outside the model. The surge wave extended from the Firth of Forth southwards: the water level in those regions was enhanced by about 0.4–0.5 m. In addition to these surge conditions, the $H_s$ of the waves at the same time was exceeding 7 m in the same areas (see Figures 14 and 15). Figures 11 and 12 show the wave field at two different times in the storm, at 00:30 and at 02:00 UTC, respectively. The swell wave effect was very low, but contributed to the enhancement of $H_s$ up to 0.5 m, while on the coastline the contribution of the WCI was very strong. At 02:00 UTC on 31 March 2010 (Fig. 15), when the storm reached the coastline, WCI increased $H_s$ by up to 2.5 m in many locations near the Firth of Forth (see Fig. 15d). Figures 14f and 15f show high $H_s$ swell waves at the entrance of the Firth of Forth. These were waves generated by the large storm shown in Figure 14e, but are no longer influenced by the local wind, but are propagat-
Figure 8. Wave conditions during the 25–26 February 2010 storm: (a) comparison between coupled and uncoupled modelled $H_s$ (m) with observed data in Firth of Forth wave gauge; (b) comparison between coupled and uncoupled modelled $H_s$ (m) in the Aberdeen wave gauge; no observation data were available from this wave gauge during this storm; (c) comparison between coupled and uncoupled modelled $T_m$ (s) with observed data in Firth of Forth wave gauge; (d) comparison between coupled and uncoupled modelled $T_m$ (s) in the Aberdeen wave gauge, no observed data were available from this wave gauge during this storm.

Figure 9. The modelled $H_s$ on the east coast of Scotland at 12:30 UTC on 26 February 2010: (a) coupled model (mean WCI included), (b) uncoupled model (mean WCI not included), (c) difference between coupled and uncoupled models, (d) difference between coupled and uncoupled models in the Moray Firth area, (e) windsea waves, (f) swell waves.

3.4.3 The 19 June 2010 storm

The third storm that is considered in this paper was one that generated high off-shore wave conditions, with swell propagating to the coastline. This is an example of how the coupling of swell and windsea waves could lead to extreme wave conditions, with significant wave height exceeding 6 m offshore and 4–5 m on the coastline. Figure 16 shows the pressure conditions between 18 and 20 June 2010. On 17 June 2010 (not shown) a system of low pressure was generated between Greenland and Iceland. This minimum moved quickly to the Scandinavian peninsula, intensifying and remaining in the area of Sweden and Norway for 72 h. This low pressure caused strong winds in the northern North Sea and consequently the generation of waves in the area between the Norway and Scotland. Recorded wave conditions in the Firth of Forth are compared with the model output (Fig. 17a–c) and model output from the Aberdeen wave rider.
Figure 10. The modelled $H_s$ on the east coast of Scotland at 19:00 UTC on 26 February 2010: (a) coupled model (mean WCI included), (b) uncoupled model (mean WCI not included), (c) difference between coupled and uncoupled models, (d) difference between coupled and uncoupled models in the Moray Firth area, (e) windsea waves, (f) swell waves.

location is shown (Fig. 17b–d). The model demonstrates the wave conditions present during this storm well (both for wave heights and periods) and the results show the limited effect of the WCI in those locations. This field of waves arrived at the Scottish coastline at the same time as the low pressure was generating high waves in the bulk of the North Sea, causing two trains of waves to be in the same place at the same time. This condition, known as crossing or bimodal sea, is quite common in the North Sea (Guedes Soares, 1984). The model hindcasted that the storm offshore was at its maximum near 16:00 UTC on 19 June 2010 (Fig. 18). At 16:00 UTC on 19 June 2010, the modelled offshore, mid-North Sea, windsea-generated waves peaked at $H_s \sim 5$ m (Fig. 14e), whereas the swell waves were a little smaller with $H_s \sim 3$–$4$ m (Fig. 18f). Further north, in the Moray Firth, the swell waves dominated with the swell having $H_s \sim 6$ m and the windsea having $H_s \sim 2$ m. The resulting predicted wave field had $H_s > 6$ m (Fig. 18b). In the Moray Firth, an $H_s$ of more than 5 m was recorded. However, at this time, the coupling between currents and waves caused a decrease of the significant wave height at the coastline (Fig. 18c). In some locations $H_s$ was reduced by more than 0.5 m (see Fig. 18c–d). Three hours later (Fig. 19), the turning tidal currents enhanced the waves by more than 1.5 m in coastal locations. In this storm, the WCIs play a role in the enhancement of the wave conditions: spatially, the effect (Figs. 18–19) is significant on the coastline. In addition, the windsea wave field is significantly enhanced by swell waves, and the bimodal sea conditions are effective in changing the $H_s$ due to the interactions between swell and windsea waves.

### 3.5 Effect of WCI on the wave spectra

Considering the second storm (30–31 March 2010) we analysed the effect of the WCI on the 1-D and 2-D spectra. Modelled spectra were extracted from the model output in three locations in correspondence with the wave gauges, and the output with and without WCIs was analysed (Figs. 20–22). Some significant variation of the energy density of the spectra ($\approx 20\%$) were seen for the considered storm, in particular for the Aberdeen wave gauge, but also for the Firth of Forth wave gauge, in which high waves were recorded; the major changes were reported near the spectral peak. The model also predicted a shift of the spectral peak and variation in swell magnitude. Since large variations were recorded for the Aberdeen wave gauge and the Firth of Forth wave gauge, we analysed the modelled directional spectra with and without WCIs for the considered storm. In Figs. 23–24, we show the results for the 2-D spectrum, in which not only the distribution of the energy with the frequency was shown but also the distribution with the angle. Variation in the magnitude of the spectral energy with the angle along with small variation in the direction of the wave train were modelled.

Similar results for the spectrum variations are reported in Rusu (2010) for the WCIs at the mouth of Danube, while similar spectral changes were identified in laboratory experiments, as in Toffoli et al. (2011b) and Toffoli et al. (2015).

### 4 Conclusions

In this study we presented a model capable of hindcasting surge and storms on the east coast of Scotland. The combination of spring tide, strong wind, and high waves can be extremely threatening in coastal areas. The North Sea is one of the areas most affected by this forcings. Storms in North Sea can generate extremely high waves as well as rogue waves (Ponce de León and Guedes Soares, 2014).

Results indicate that WCIs play a fundamental role in the wave propagation during severe storms in the coastal areas, while for the open sea, the maximum contribution of this interaction is less than 0.5 m of magnitude. The results are consistent with other studies of WCI in other parts of the

world, such as in the southern North Sea (Osuna and Monbaliu, 2004), in which the difference based on the normalized rms difference of a 1-month period is about 3% for the $H_s$ and has an rms of 20% for $T_m$ (Tables 3 and 4 of Osuna and Monbaliu, 2004). Such is the case in the northern Adriatic Sea in the shallow areas between the Venetian Lagoon, the Gulf of Trieste, and the Istrian peninsula, where deviations up to 1 m were modelled during bora and sirocco conditions (Benetazzo et al., 2013).

The validation shows that the model performs reasonably well during both calm periods and storms for waves, and also performs well for tides and surges.

During severe storms, in particular when the low pressure was over England and Scotland, it was found that the WCIs are significant, causing an increase or decrease in $H_s$ that can exceed 2 m in some coastal areas, depending on the direction of the wave field compared to the current. A similar result was found for the peak spectral wave period: Fig. 4 shows that in the time period considered here the largest deviation of wave periods due to WCI is in the estuarine areas of the east coast, with rms deviations of more than 1.2 s.

Wave propagation in the Firth of Forth during storms generated in the mid-North Sea is driven by trains of swell waves detaching from the open sea storm. During the stormy periods considered here, the windsea waves in the Firth of Forth did not exceed 3.5 m in the outer area of the estuary and 1 m in the inner part, while the swell field exceeded 5 m at the entrance of the Firth of Forth. In the inner firth, the swell waves have a similar magnitude to the windsea waves. Conversely, the area of the estuary of the Firth of Forth is mainly driven by locally generated waves. A similar behaviour was noticed in the other two estuarine areas on the east coast: the Tay estuary and the Moray Firth.

The northeast coast of Scotland is more exposed to swell arriving from the North Atlantic and the Norwegian Sea, while the central and southern parts are more exposed to local windsea waves and to storms generated in the bulk of the North Sea.

Spectra were also considered in the analysis of the WCIs: spectral variations, in particular in the energy peak, were significant and exceeding 20% in some cases. Wave periods are adequately modelled by the model presented in this paper. Wave models, however, have a large error for the wave period, since they do not include nonlinear quasi-resonant interactions (Onorato et al., 2002; Janssen, 2003; Waseda et al., 2009) that are also fundamental for the correct estimation of the $H_s$ when the spectrum is narrow. In addition, wave periods from satellite data are often very difficult to estimate (Gommenginger et al., 2003). Another limitation of the study is that no surge boundaries were available, so the water level
Figure 12. Modelled currents and waves conditions in the Aberdeen wave gauge location during the 30–31 March 2010 storm (depth of the mooring location is 10 m).

and the current fields were only due by tides and the local field of wind and pressure. This led to an overall underestimation of the strength of the current and a possible underestimation of the total effect of the WCI.

The model also has forecasting capabilities, in particular when nested with large-scale models, such as the North Atlantic model (Venugopal and Nemalidinne, 2014, 2015). A limitation of the model is that the MIKE by DHI software does not allow an online coupling between waves and tides, slowing the simulation process. In fact, currents and waves are simulated by different modules and it is not possible to perform a direct coupling. For this work, the currents were simulated first and then the output data were saved in order to use them as input for the wave model. Another limitation of the model, due to the one-way coupling, is that we cannot study the effect of the wave setup and setdown on the surge water level (Longuet-Higgins and Stewart, 1962; Bowen et al., 1968), and most importantly, the wave radiation effect on the current field itself (Bennis et al., 2011; Michaud et al., 2011). Previous work on this interaction shows that the modification of the current field is more important in very shallow water areas (< 10 m depth). In this paper, however, we were more interested in the effect of the current field on the wave. Future work will focus on understanding what effect the waves have on the current dynamics on the east coast of Scotland. This will be implemented by first running the wave model, then using the wave radiation in the hydrodynamic model to estimate the enhancement of the water level due to waves near the shoreline and to estimate the variation of the current due to the wave radiation stress.

This research also underlines the importance of high-resolution regional-scale models for the understanding of sea dynamics and the forecasting of dangerous sea states: larger models usually have inadequate resolution to estimate the effect of such processes near the coastline. Future work will be focused on the hindcasting of freakish wave states based on the estimation of the kurtosis from the parameters of the model (Janssen, 2003; Tamura et al., 2009; Ponce de León and Guedes Soares, 2014) and on the sediments resuspension in the area of Stonehaven (Heath et al., 2015), which is an intensive study site for suspended sediment and other biological variables in the water column (Serpetti et al., 2011, 2012).

Figure 13. The modelled surge wave due to the local wind and pressure at 02:00 UTC on 31 March 2010.
Figure 14. The modelled $H_s$ on the east coast of Scotland at 00:30 UTC on 31 March 2010: (a) coupled model (mean WCI included), (b) uncoupled model (mean WCI not included), (c) difference between coupled and uncoupled models, (d) difference between coupled and uncoupled models in the Firth of Forth area, (e) windsea waves, (f) swell waves.

Figure 15. The modelled $H_s$ on the east coast of Scotland at 02:00 UTC on 31 March 2010: (a) coupled model (mean WCI included), (b) uncoupled model (mean WCI not included), (c) difference between coupled and uncoupled models, (d) difference between coupled and uncoupled models in the Firth of Forth area, (e) windsea waves, (f) swell waves.
Figure 16. The mean sea level pressure fields (hPa) before and during the 19 June 2010 storm.
Figure 17. Wave conditions during the 19 June 2010 storm: (a) comparison between coupled and uncoupled modelled $H_s$ (m) with observed data in Firth of Forth wave gauge; (b) comparison between coupled and uncoupled modelled $H_s$ (m) in the Aberdeen wave gauge, no observed data were available from this wave gauge during this storm; (c) comparison between coupled and uncoupled modelled $T_m$ (s) with observed data in Firth of Forth wave gauge; (d) comparison between coupled and uncoupled modelled $T_m$ (s) in the Aberdeen wave gauge, no observed data were available from this wave gauge during this storm.

Figure 18. The modelled $H_s$ on the east coast of Scotland at 16:00 UTC on 19 June 2010: (a) coupled model (mean WCI included), (b) uncoupled model (mean WCI not included), (c) difference between coupled and uncoupled models, (d) difference between coupled and uncoupled models in the Firth of Forth area, (e) windsea waves, (f) swell waves.
Figure 19. The modelled $H_s$ on the east coast of Scotland at 19:00 UTC on 19 June 2010: (a) coupled model (mean WCI included), (b) uncoupled model (mean WCI not included), (c) difference between coupled and uncoupled models, (d) difference between coupled and uncoupled models in the Firth of Forth area, (e) windsea waves, (f) swell waves.

Figure 20. Modelled 1-D spectrum in the Firth of Forth wave gauge, the red line is the coupled model (with WCIs incorporated), while the blue line is the uncoupled model: (a) 31 March 2010 at 00:30, (b) 31 March 2010 at 01:15, (c) 31 March 2010 at 02:00, (d) 31 March 2010 at 04:15, (e) 31 March 2010 at 06:00, and (f) 31 March 2010 at 08:30 UTC.
Figure 21. Modelled 1-D spectrum in the Moray Firth wave gauge, the red line is the coupled model (with WCIs incorporated), while the blue line is the uncoupled model: (a) 31 March 2010 at 00:30, (b) 31 March 2010 at 01:15, (c) 31 March 2010 at 02:00, (d) 31 March 2010 at 04:15, (e) 31 March 2010 at 06:00, and (f) 31 March 2010 at 08:30 UTC.

Figure 22. Modelled 1-D spectrum in the Aberdeen wave gauge, the red line is the coupled model (with WCIs incorporated), while the blue line is the uncoupled model: (a) 31 March 2010 at 00:30, (b) 31 March 2010 at 01:15, (c) 31 March 2010 at 02:00, (d) 31 March 2010 at 04:15, (e) 31 March 2010 at 06:00, and (f) 31 March 2010 at 08:30 UTC.
Figure 23. Polar plot of the modelled 2-D directional spectrum (energy density, \(m^2 \times s/\)degrees) in the Firth of Forth wave gauge, red indicates the contour plot of the coupled model spectrum (with WCIs incorporated), while black indicates the contour plot of the uncoupled model. Contour lines are plotted every 0.01 \(m^2 \times s\) degrees\(^{-1}\): (a) 31 March 2010 at 00:30, (b) 31 March 2010 at 01:15, (c) 31 March 2010 at 02:00, (d) 31 March 2010 at 04:15, (e) 31 March 2010 at 06:00, and (f) 31 March 2010 at 08:30 UTC.

Figure 24. Polar plot of the modelled 2-D directional spectrum (energy density, \(m^2 \times s/\)degrees) in the Aberdeen wave gauge, red indicates the contour plot of the coupled model spectrum (with WCIs incorporated), while black indicates the contour plot of the uncoupled model. Contour lines are plotted every 0.01 \(m^2 \times s\) degrees\(^{-1}\): (a) 31 March 2010 at 00:30, (b) 31 March 2010 at 01:15, (c) 31 March 2010 at 02:00, (d) 31 March 2010 at 04:15, (e) 31 March 2010 at 06:00, and (f) 31 March 2010 at 08:30 UTC.
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